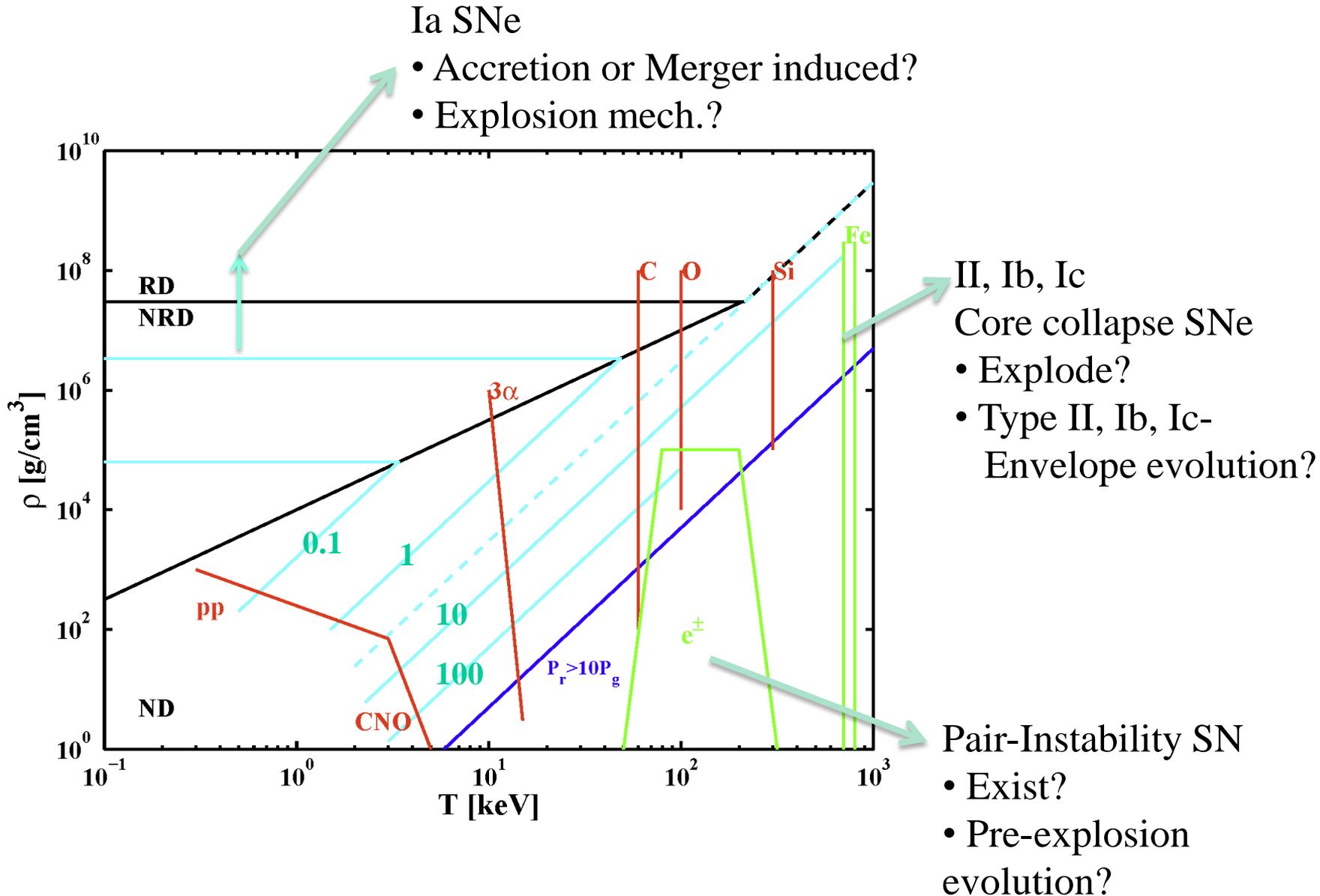


# A New Mechanism for Deflagration to Detonation Transitions

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# Supernova (SN) types & Open Q's



# Ia's

## □ The common view

- Ignition of C/O under degenerate conditions (White Dwarf) → thermonuclear "runaway" → No NS remnant,  $\sim 10^{51}$  erg (1MeV/nucleon= $2 \times 10^{51}$  erg/ $M_{\text{sun}}$ )
- Support: Composition,  $v$ , Light curves (E,  $v$ )

[Hoyle & Fowler 60]

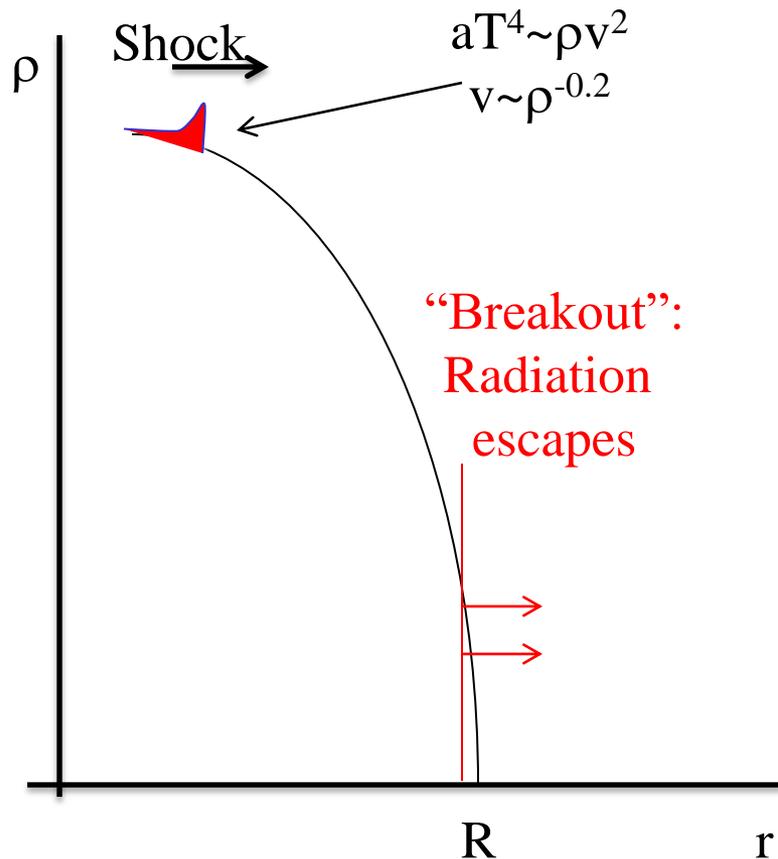
## □ Open Q's

- Progenitors: Accretion/merger?
- How is the explosion triggered?
- How does the burning propagate: Pure deflagration or DDT?
- DDT mechanism?
- Why so uniform?

[eg Hillebrandt & Niemeyer 00]

# A recent development

- Wide field surveys (eg PTF): post-Breakout-detection

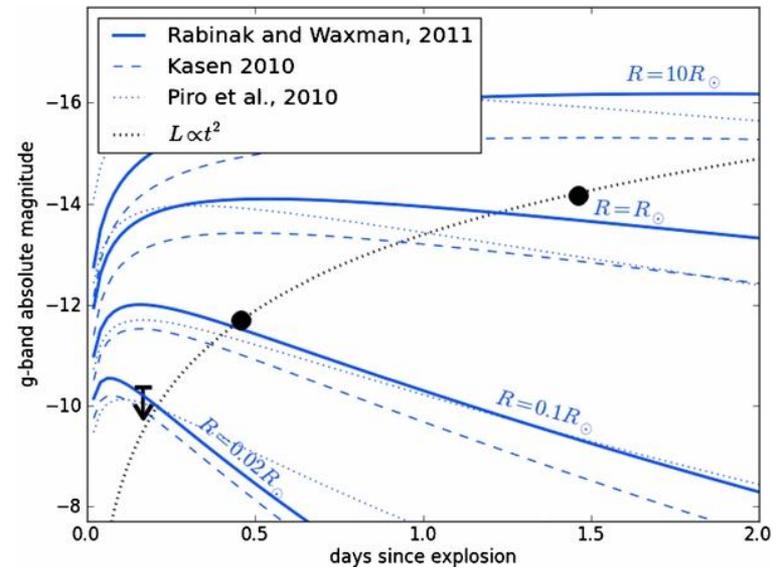


- Ia SN 2011fe early non-detection:

First direct determination of compact prog. Radius:

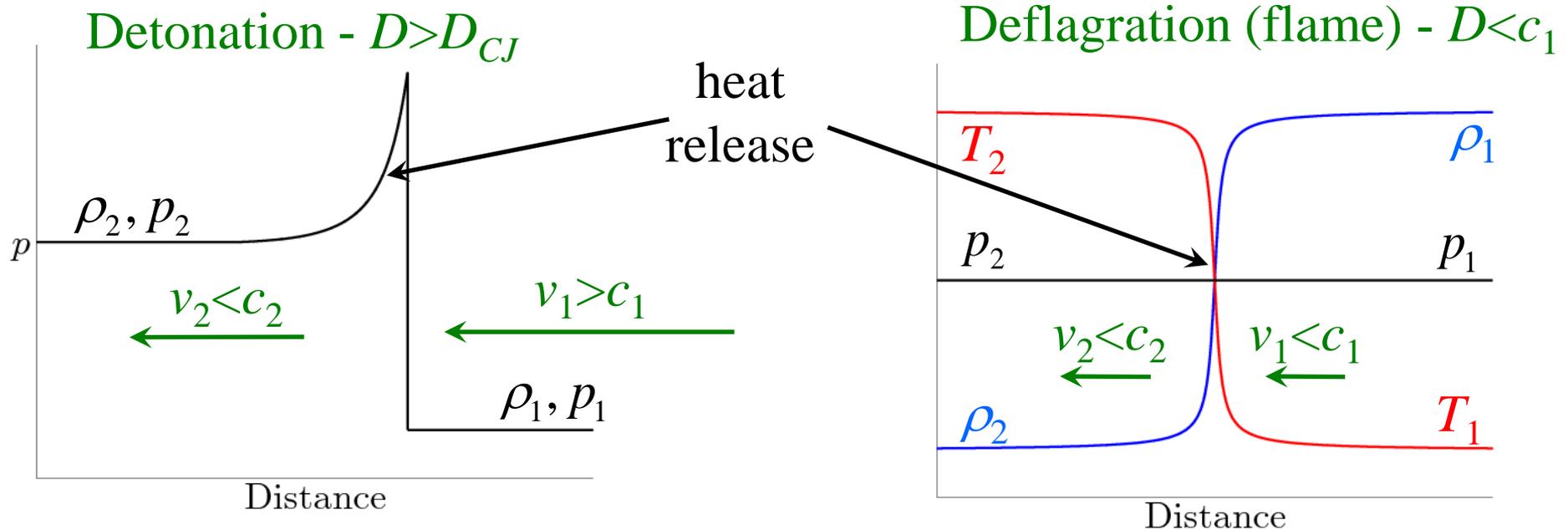
$$R < 10^{10} \text{ cm [Nugent et al 11]}$$

$$R < 10^9 \text{ cm [Bloom et al 11]}$$



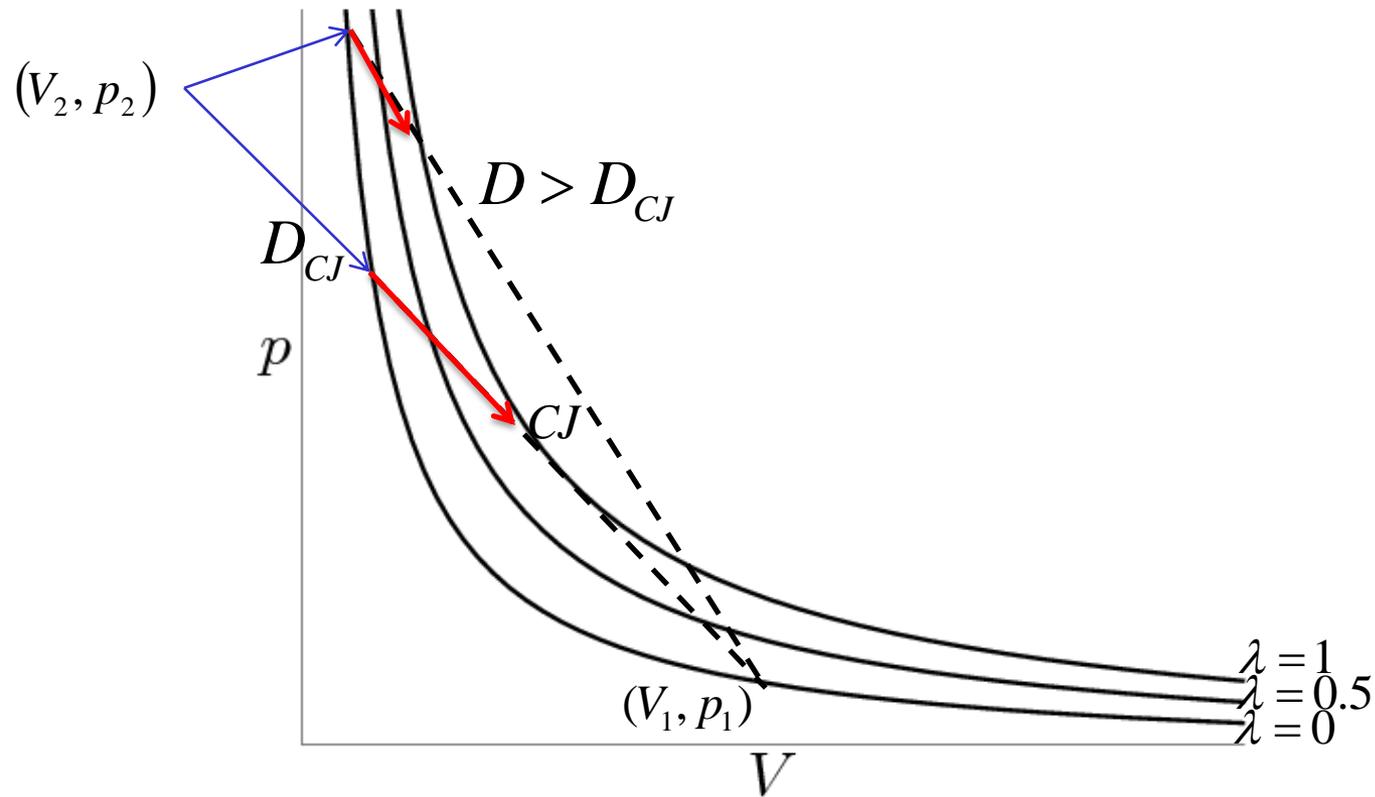
[Bloom et al. 11]

# Ia: Thermonuclear combustion of C/O WD



# ZND Theory

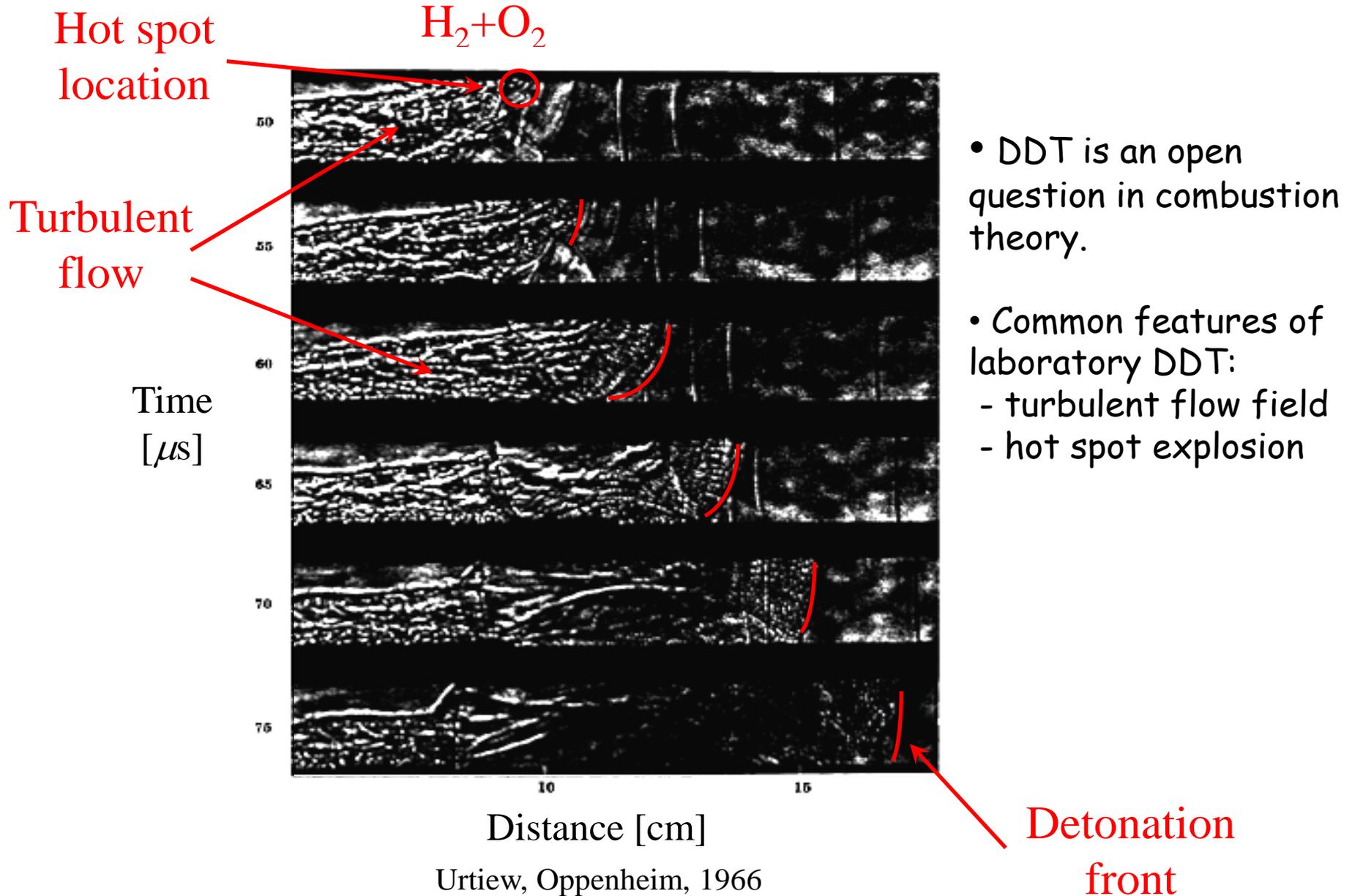
shock  $\Rightarrow$  ignition  $\Rightarrow$  reaction zone



# Ia DDT

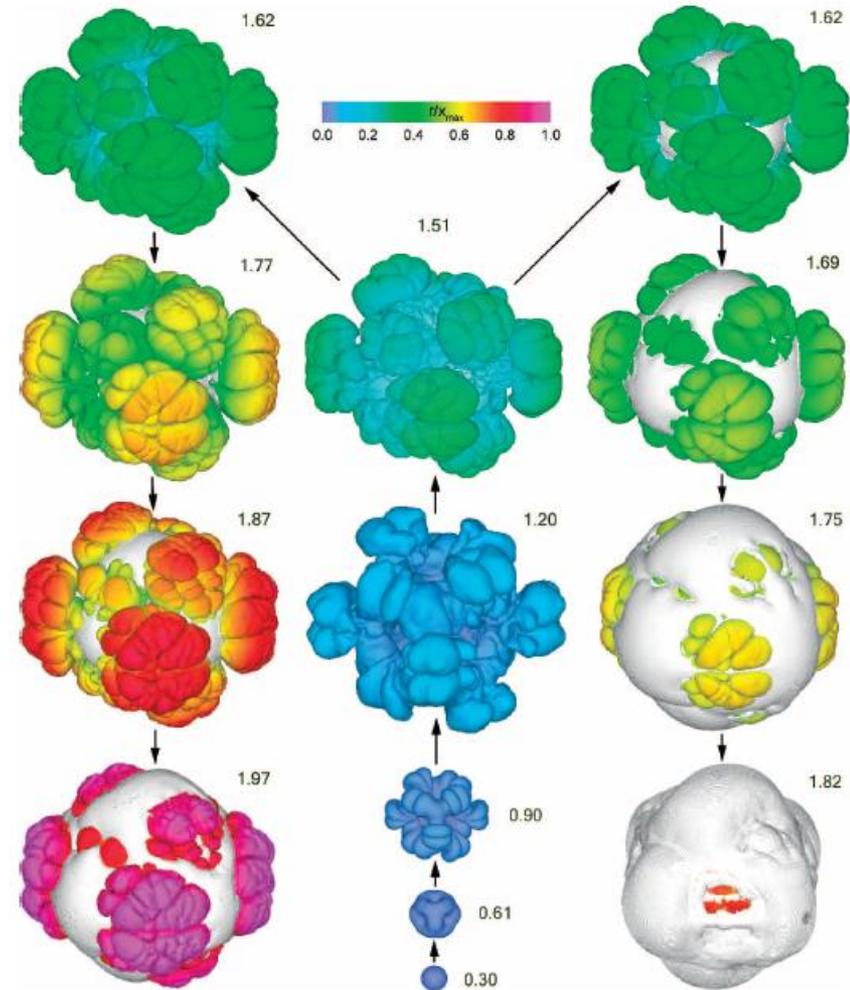
- Ejecta composition suggests  
Deflagration-to-Detonation-Transition (DDT)
- Pure detonation- complete burning to Fe group
- Sub-sonic deflagration: expansion → lower density burning → intermediate mass elements
- Pure deflagration difficulties:  $v$ , mixed composition

# DDT in the laboratory



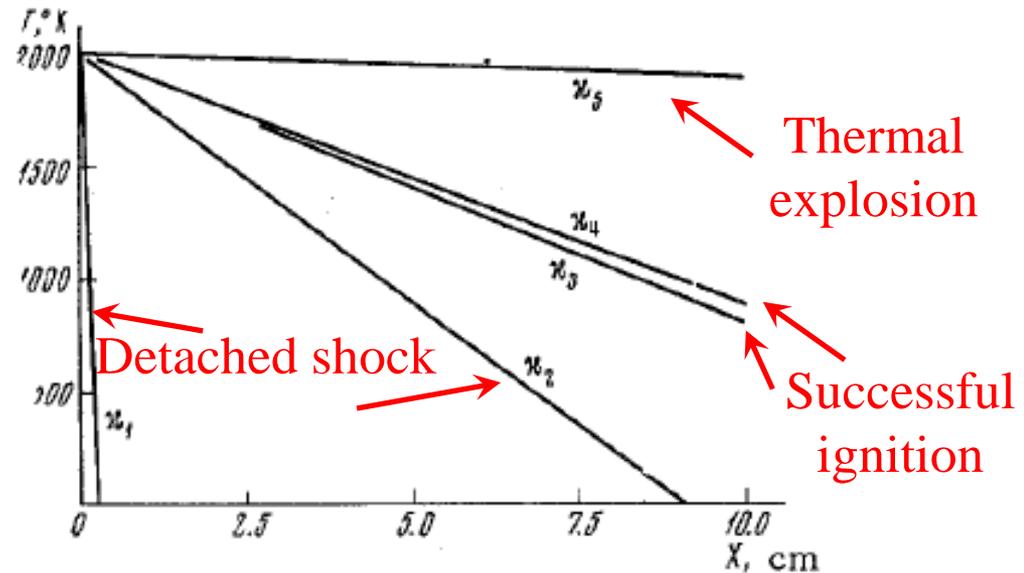
# DDT in Ia explosions

- Rayleigh-Taylor instability  
→ turbulent flame  
→ burned/unburned material mix
- Detonation initiated by artificially introducing a Hot spot



# DDT mechanisms: I

## Initial temperature distributions



- Zel'dovich et al., 1970  
A spatial gradient in chemical induction time  
→ A spontaneous reaction wave with  $v_p \approx D_{CJ}$ .

?? How gradients are maintained over the critical length?  
( $10^4$ – $10^5$  cm @  $10^7$  g/cc, Seitenzahl et al. 09)

# DDT mechanisms: II

- Our Mechanism:

- Converging shock waves ignite detonation, provided the radius at which they become strong exceeds  $R_{\text{crit}}$  (may be tested experimentally).

- DDT may be due to converging shocks produced by the turbulent deflagration flow, which reaches sub (but near) sonic velocities on scales  $\gg R_{\text{crit}}$ .

?? Are such converging shocks indeed produced in deflagration flows?

# A simple model

- EOS:  $\varepsilon(p, V, \lambda) = \frac{pV}{\gamma - 1} - \lambda Q$
- Burning:  $\frac{d\lambda}{dt} = \kappa \left( \frac{\rho}{\rho_0} \right)^n (1 - \lambda)^m e^{-(p_A/\rho_0)/(p/\rho)}$

- Initial conditions: Converging shock

$$\dot{R} = -D_{CJ} \left( \frac{R}{R_{CJ}} \right)^\delta, \quad D_{CJ} = \sqrt{2Q(\gamma^2 - 1)}$$

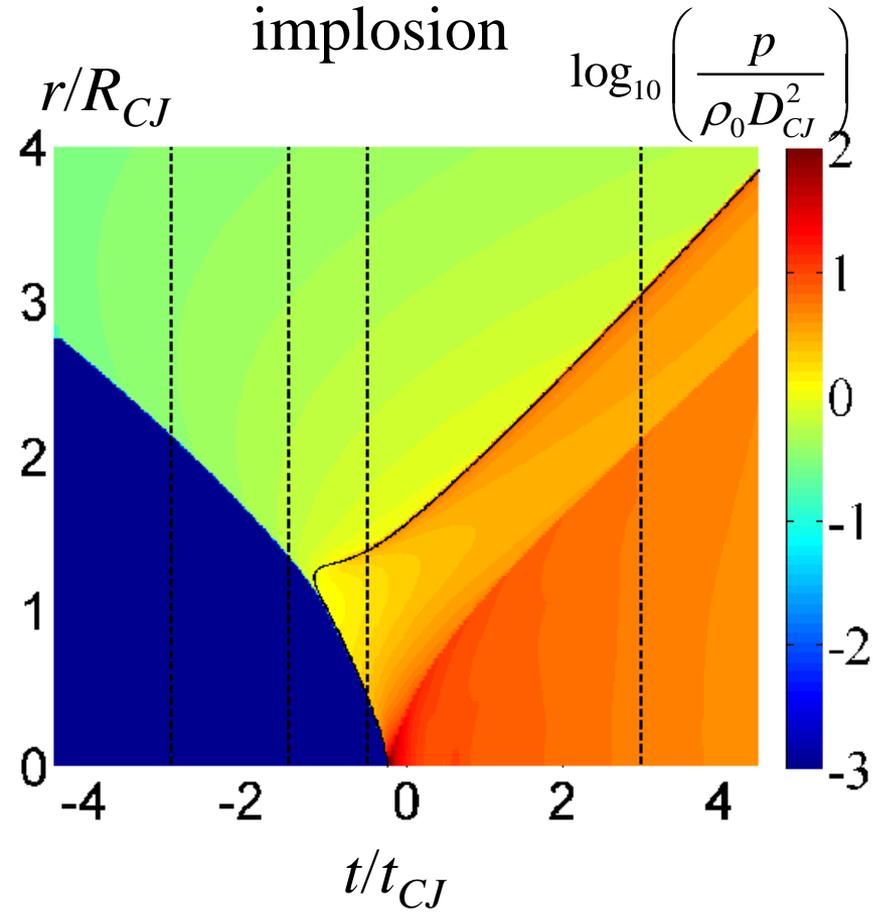
- 5 dimensional parameters ( $\rho_0, D_{CJ}, R_{CJ}, \kappa, p_A$ )  
→ fully determined by 2 dimensionless parameters:

$$\tau \equiv \frac{p_A}{\rho_0 D_{CJ}^2}, \quad \text{and } \theta \text{ (TBD).}$$

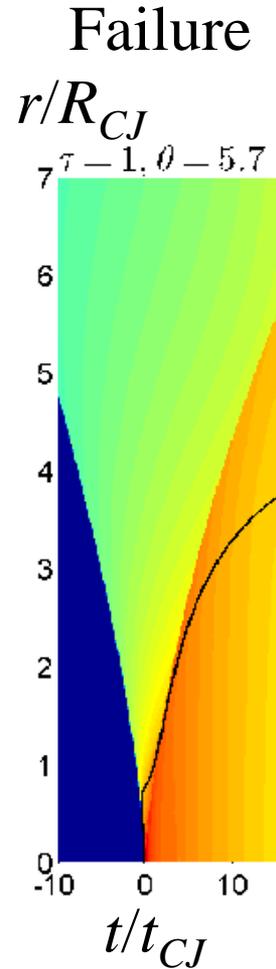
(+  $\gamma, n, m$ ).

# Some numerical examples

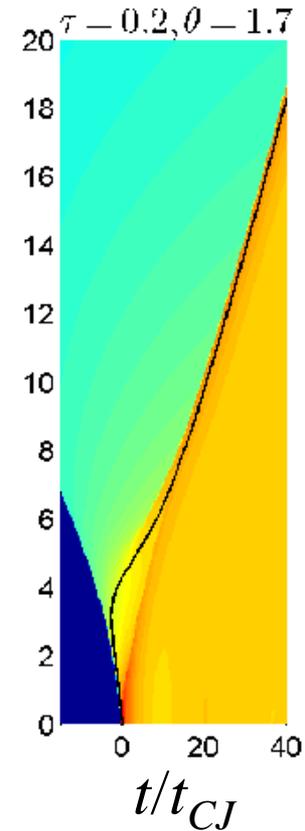
Ignition by  
implosion



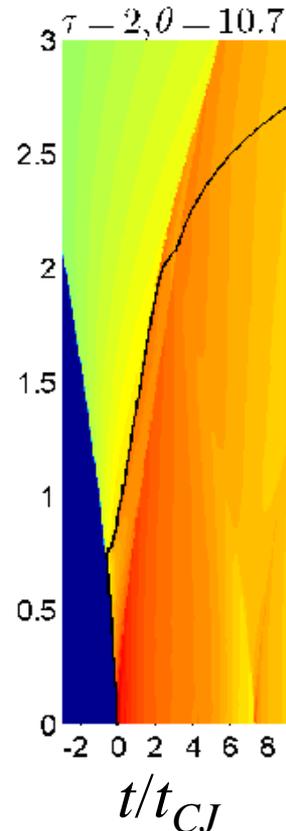
Failure



Ignition by  
explosion



Failure



# Approximate analytic results

- Ignition criterion:

$$\frac{Q}{t_q} \min[t_q, t_h] > f Q_h = f \frac{2}{(\gamma+1)^2} \dot{R}^2 \quad \text{for some } R,$$

$$\text{where } t_q^{-1} \equiv \frac{d \ln \lambda}{dt}, t_h^{-1} \equiv \frac{d \ln R}{dt}.$$

- Translates to:  $\theta > f g(\xi)$  for some  $\xi > \left[ \frac{8(\gamma-1)}{\gamma+1} \right]^{-1/2\delta}$ ,

$$\text{where } \xi \equiv R / R_{CJ},$$

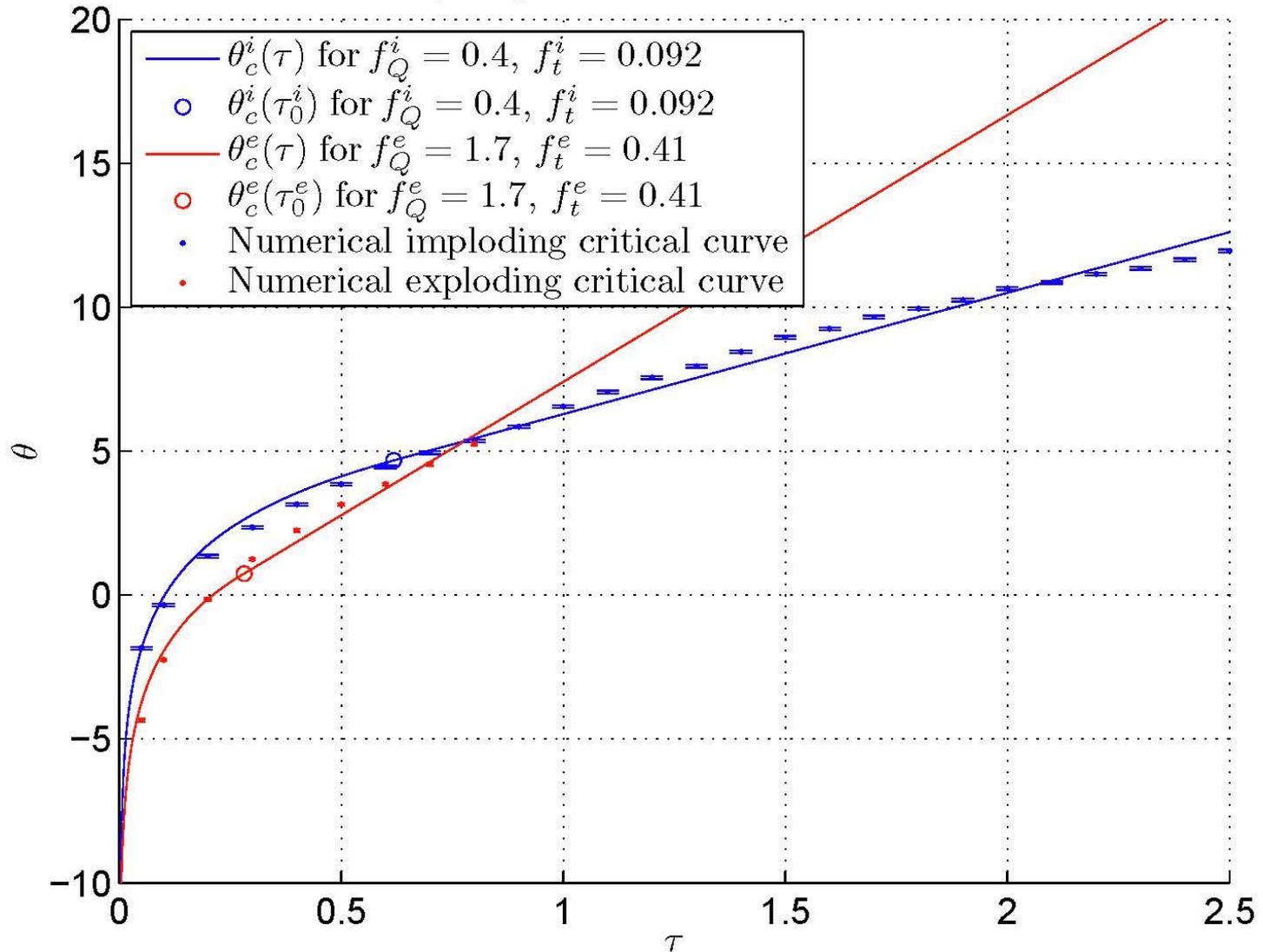
$$g(\xi) = \eta \xi^{-2\delta} - (1-3\delta) \ln \xi,$$

$$\theta \equiv \ln \left( \frac{\kappa R_{CJ}}{(1-\delta) D_{CJ}} \right) + n \ln \left( \frac{\gamma+1}{\gamma-1} \right) + 2\delta \ln \left( \frac{4(\gamma-1)}{\gamma+1} \right).$$

→  $\theta > \theta_c(\tau)$ , i.e.  $R_{CJ}$  should exceed a critical value.

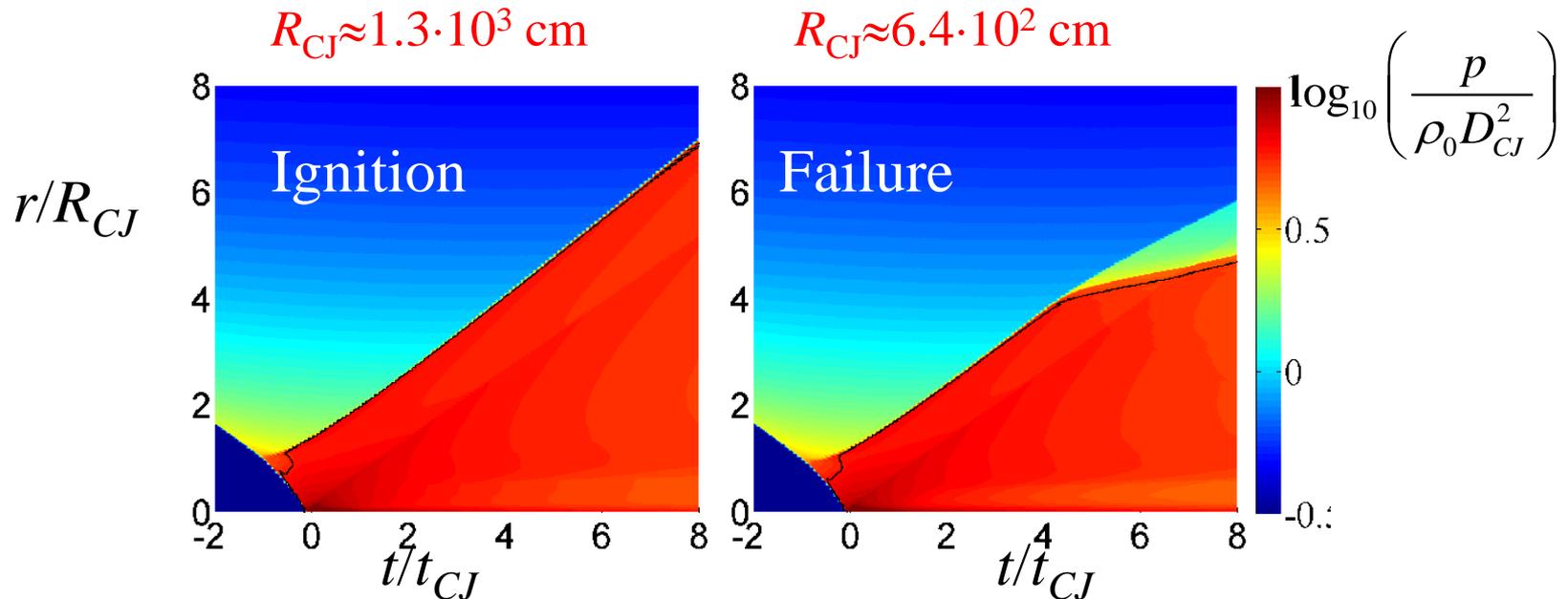
# Critical radii: Analytic vs. Numeric

$$\gamma = 5/3, m = 1, n = 0, \nu = 3$$



# Critical radii: Lab & Ia's

- Laboratory stoichiometric acetylene-air experiments:
  - $R_{CJ} \sim 5\mu$ ,  $R_{M=2} \sim 100\mu$ .
  - $R_{CJ}$  too small to be resolved (numerically),  $R_{M=2} \ll$  channel diameter
- Ia's: Velocity fluctuations  $\simeq 10^8$  cm/s ( $M \approx 1.2$ ) on 10 km scale,  
 $\rho_0 \simeq 10^7$  g/cc
  - Preliminary:  $R_{CJ}$  too small to be resolved,  $R_{M=2} \sim 0.1$  km  $\ll$  10 km



# Summary

- Converging shock waves ignite a detonation provided the radius at which the shock becomes strong exceeds  $R_{\text{crit}}$ .
- $R_{\text{crit}} \approx 1$  mm for typical acetylene-air experiments and  $R_{\text{crit}} \approx 10^4$  cm for the pre-detonation phase of WD in delayed-detonation scenarios of SNIa explosions.
- We suggest that the DDT observed/inferred in these systems may be due to converging shocks produced by the turbulent deflagration flow, which reaches sub (but near) sonic velocities on scales  $\gg R_{\text{crit}}$ .
- Under progress: Evolution of multidimensional perturbations during shock implosion does not suppress the ignition of detonation;  
Ia's: Realistic EOS & nuclear reaction network.
- In order to determine whether our suggested mechanism is indeed responsible for DDT, a detailed analysis of the turbulent flow is required.