

Topological transitions in electronic spectra of vortex clusters in type-II superconductors

A.S.Mel'nikov

Institute for Physics of Microstructures RAS, Nyzhny Novgorod

A.V.Samokhvalov and M.A.Silaev

◆ Introduction

Types of topological electronic transitions in the band spectra of crystals.

Fermi surface in vortex state of superconductors.

◆ Examples of topological electronic transitions in vortex matter.

de Haas – van Alfen oscillations in vortex state.

Magnetic breakdown criterion.

Quasiparticle tunneling between vortices and reconnection of Fermi surface segments.

◆ Opening of Fermi surface segments in vortex cores.

**Creation of vortex-antivortex pairs,
vortex near the surface or defect.**

Landau criterion for vortex depinning.

Two step scenario of vortex escape from the defect.

Topological electronic transitions (Lifshits-1960)

Void formation/disappearance



Quasimomentum space.

$$E(\vec{k}) = E_F$$

Topological electronic transitions (Lifshits-1960)

Void formation/disappearance

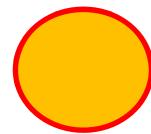
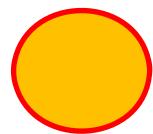


Quasimomentum space.

$$E(\vec{k}) = E_F$$

Topological electronic transitions (Lifshits-1960)

Void formation/disappearance

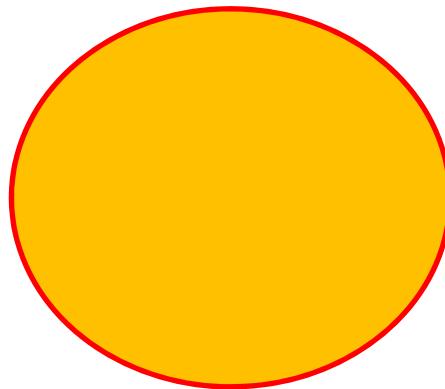
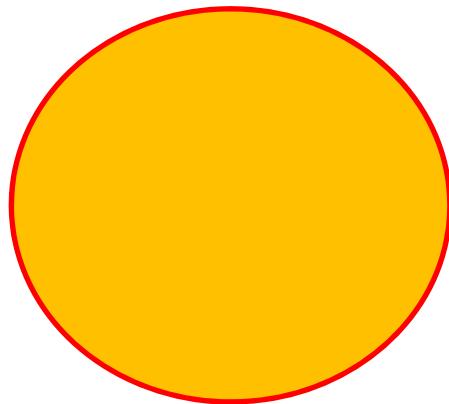


Quasimomentum space.

$$E(\vec{k}) = E_F$$

Topological electronic transitions (Lifshits-1960)

Void formation/disappearance

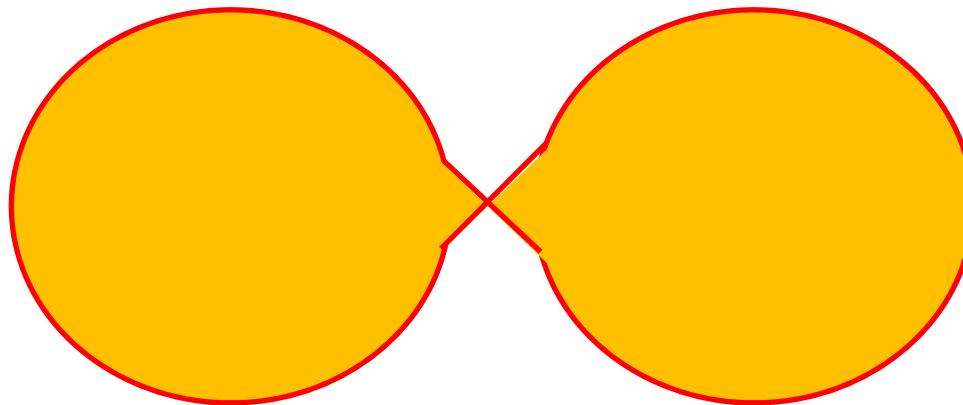


Quasimomentum space.

$$E(\vec{k}) = E_F$$

Topological electronic transitions (Lifshits-1960)

Neck formation/disruption

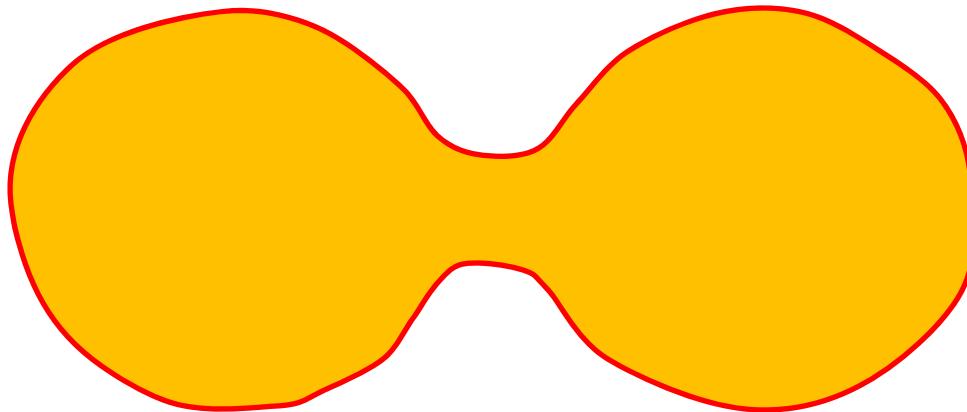


Quasimomentum space.

$$E(\vec{k}) = E_F$$

Topological electronic transitions (Lifshits-1960)

Neck formation/disruption



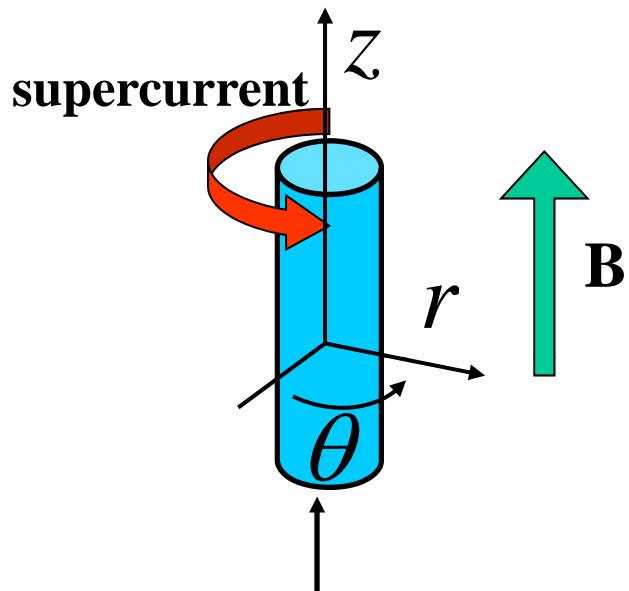
Experiment: peculiarities of DOS and transport characteristics

Quasimomentum space.

$$E(\vec{k}) = E_F$$

Fermi surface in vortex state of superconductors

Vortex line



$$\Delta = |\Delta(r)| e^{i\theta}$$

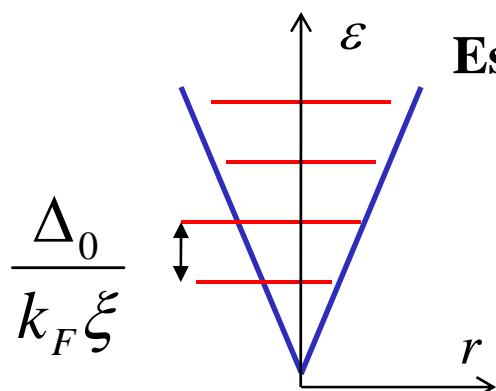
phenomenological theory of the mixed state

$$-\xi^2(T) \left(\nabla + \frac{2\pi i}{\Phi_0} \vec{A}(\vec{r}) \right)^2 \Psi = \Psi - |\Psi|^2 \Psi$$

ξ - Coherence length (core radius)

Bound fermionic states in vortex core

Superconducting gap profile:
potential well for electrons



Estimate of minigap in excitation spectrum

$$\epsilon_{\min} \sim \frac{\hbar^2}{m\xi^2} \sim \frac{\hbar^2 \Delta_0}{m\hbar\nu_F \xi} \sim \frac{\Delta_0}{k_F \xi}$$

STM vortex images

PRL 101, 166407 (2008)

PHYSICAL REVIEW LETTERS

week ending
17 OCTOBER 2008

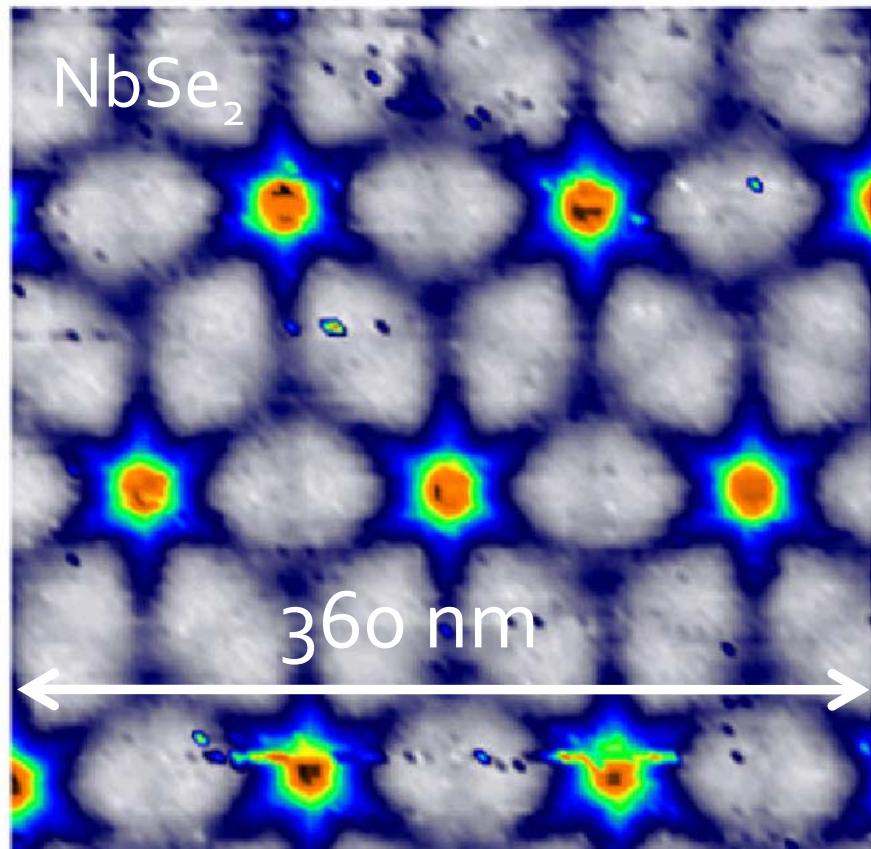
VOLUME 75, NUMBER 14

PHYSICAL REVIEW LETTERS

2 OCTOBER 1995

Superconducting Density of States and Vortex Cores of 2H-NbS₂

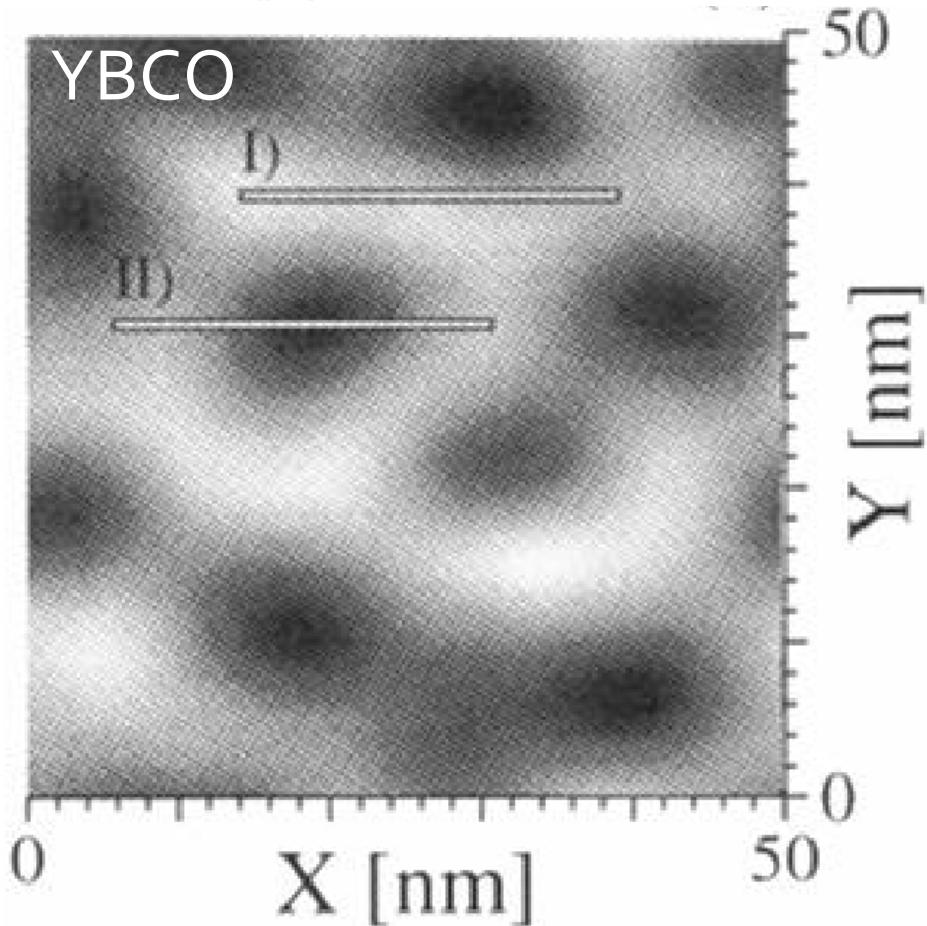
I. Guillamón,¹ H. Suderow,¹ S. Vieira,¹ L. Cario,² P. Diener,³ and P. Rodière³



PRL, 101, 166407 (2008)

Direct Vortex Lattice Imaging and Tunneling Spectroscopy of Flux Lines on YBa₂Cu₃O_{7-δ}

I. Maggio-Aprile, Ch. Renner, A. Erb, E. Walker, and Ø. Fischer



PRL, 75, 2754 (1995)

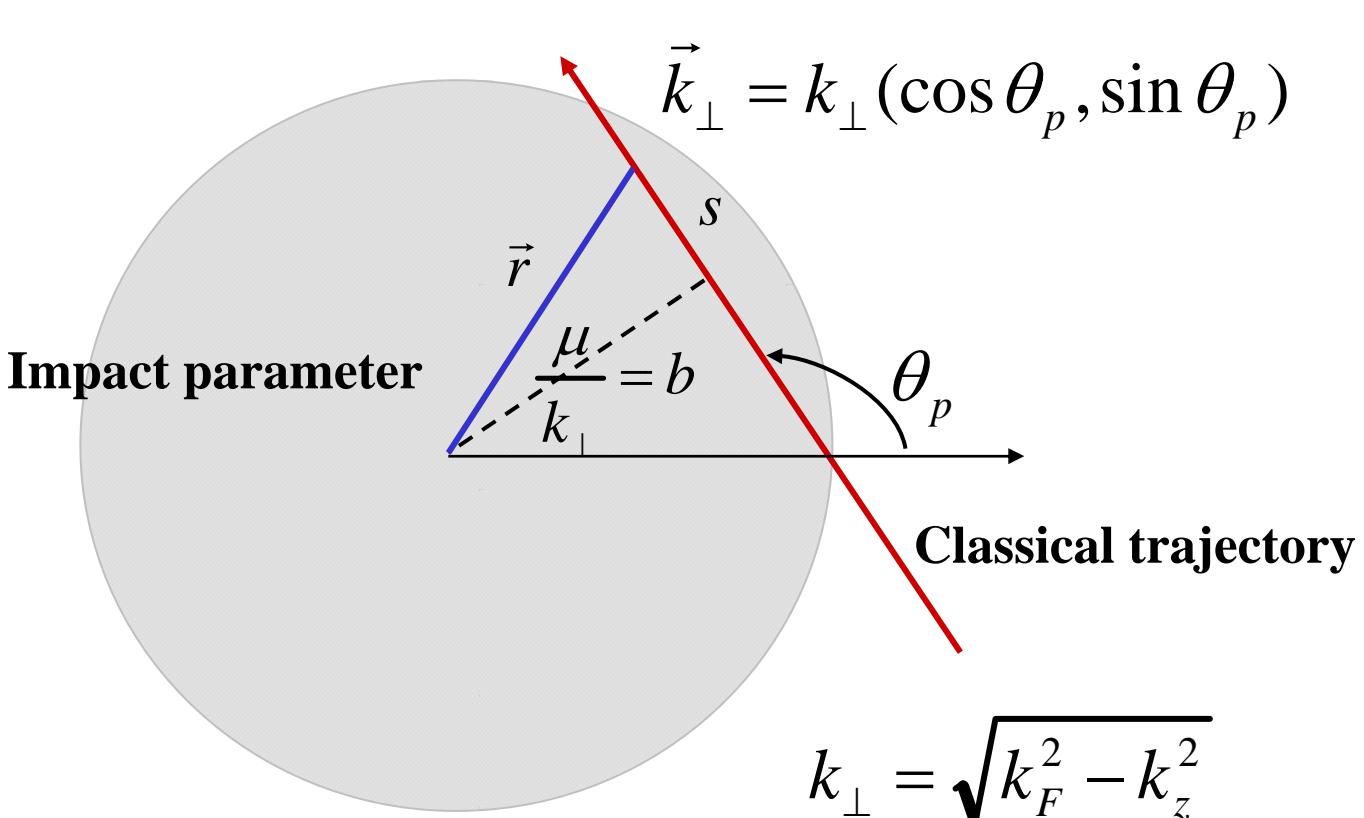
Fermi surface for electrons inside the core

Convenient parametrization of a classical trajectory:

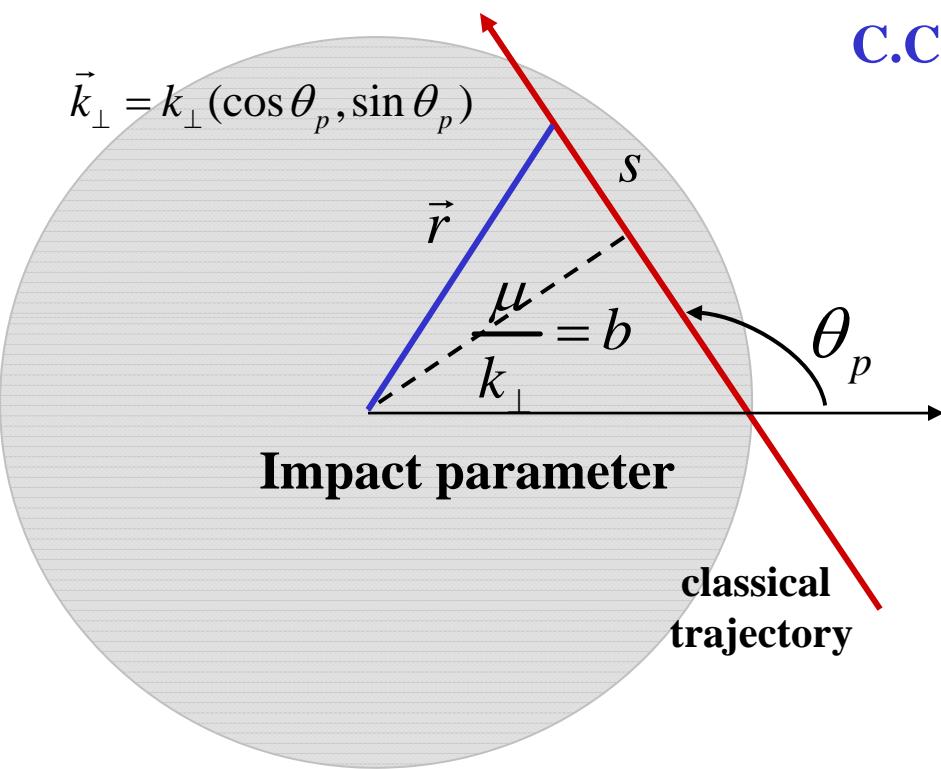
1. Orbital momentum

$$\mu = [\vec{r}, \vec{k}_\perp] \vec{z}_0 = k_\perp r \sin(\theta_p - \theta)$$

2. Trajectory orientation angle



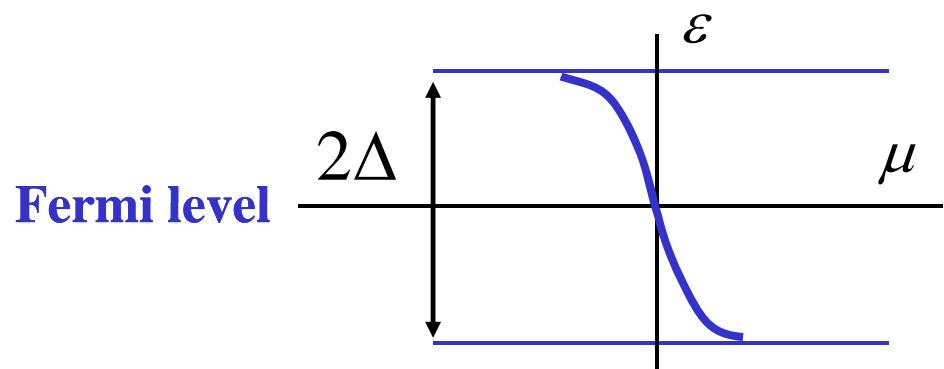
Bound quasiparticle states.



$$k_\perp = \sqrt{k_F^2 - k_z^2}$$

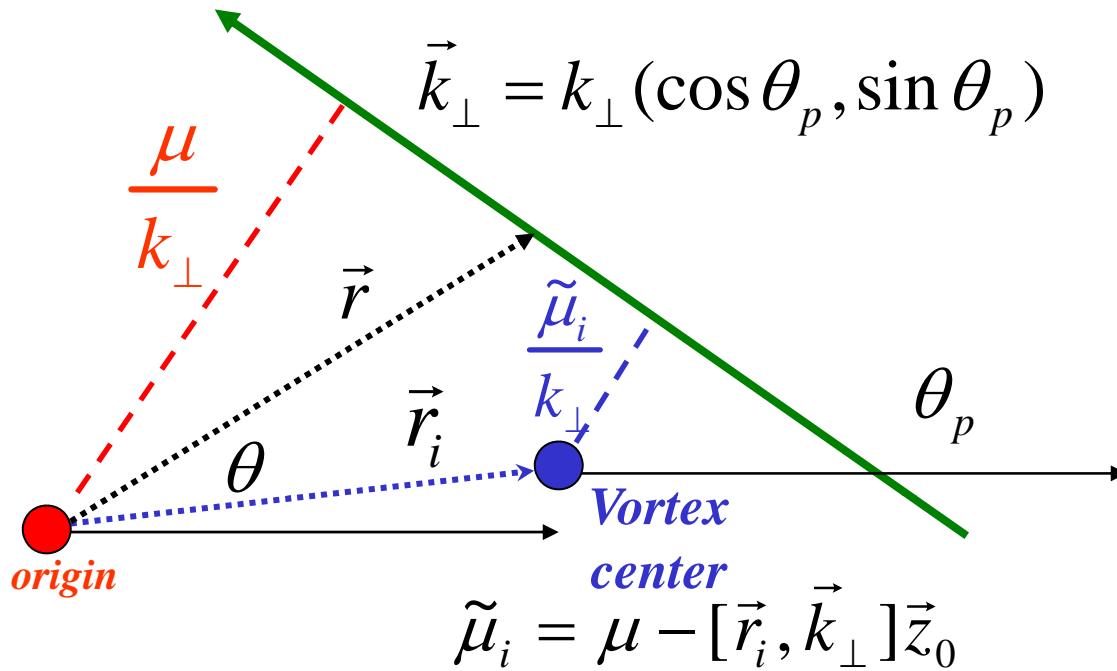
C.Caroli, P.G.de Gennes, J.Matricon (1964)

Anomalous spectral branch.



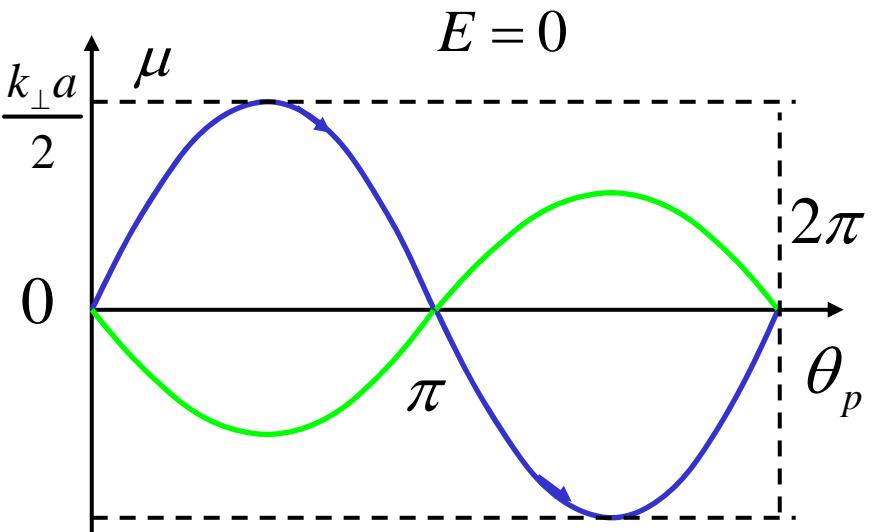
$$\epsilon_\mu(k_\perp) = -\omega\mu \approx -\frac{\mu\Delta_0}{k_\perp\xi}$$

Fermi surface for electrons inside the core



$$\varepsilon_i = -\omega \tilde{\mu}_i = -\omega \mu + \omega [\vec{r}_i, \vec{k}_\perp] \vec{z}_0$$

$$\mu_i(\theta_p) = -E/\omega + \omega [\vec{r}_i, \vec{k}_\perp] \vec{z}_0$$



Topological electronic transitions in vortex state?

Opening of Fermi surface segments = creation of vortex-antivortex pair

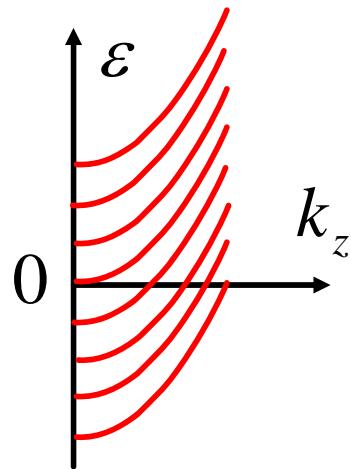
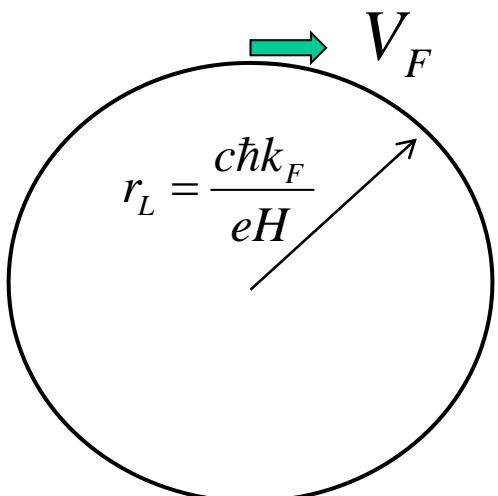
Merging and reconnection of Fermi surface segments = phase transitions in vortex matter governed by magnetic field and transport current

G.E.Volovik, JETP Lett. (1989):

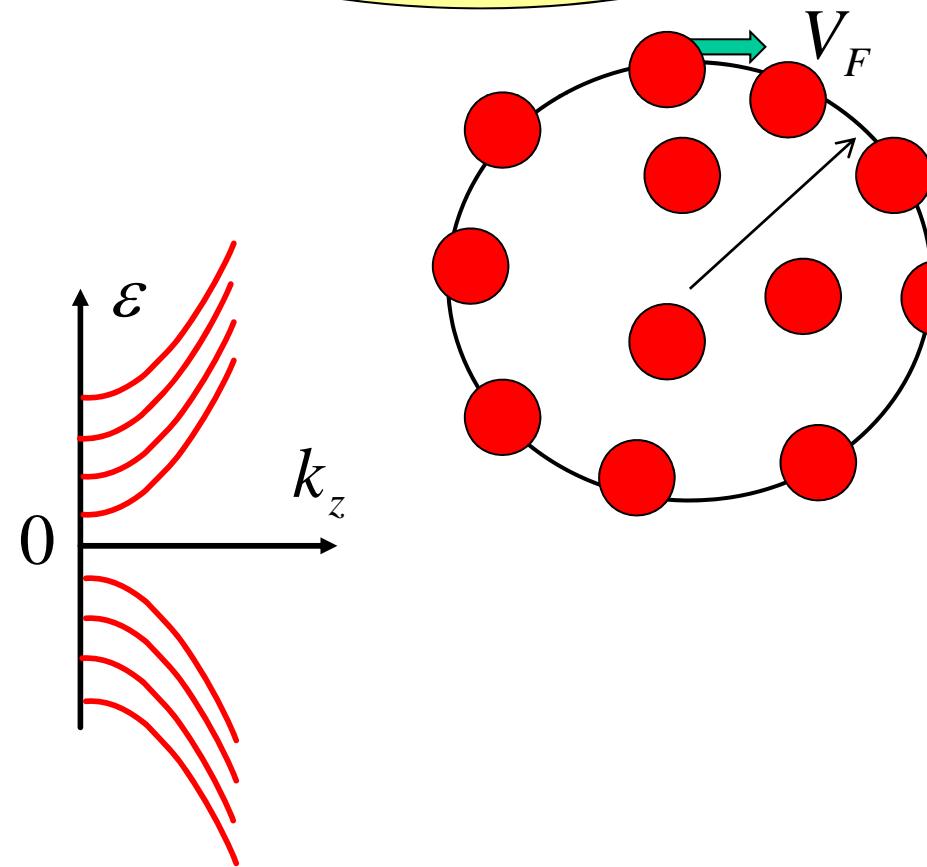
Lifshits-type transitions between vortex states with different number of zero energy modes

Simple example of topological transitions: de Haas-van Alfen oscillations.

Normal metal.
Landau spectrum.
Cyclotron orbits.

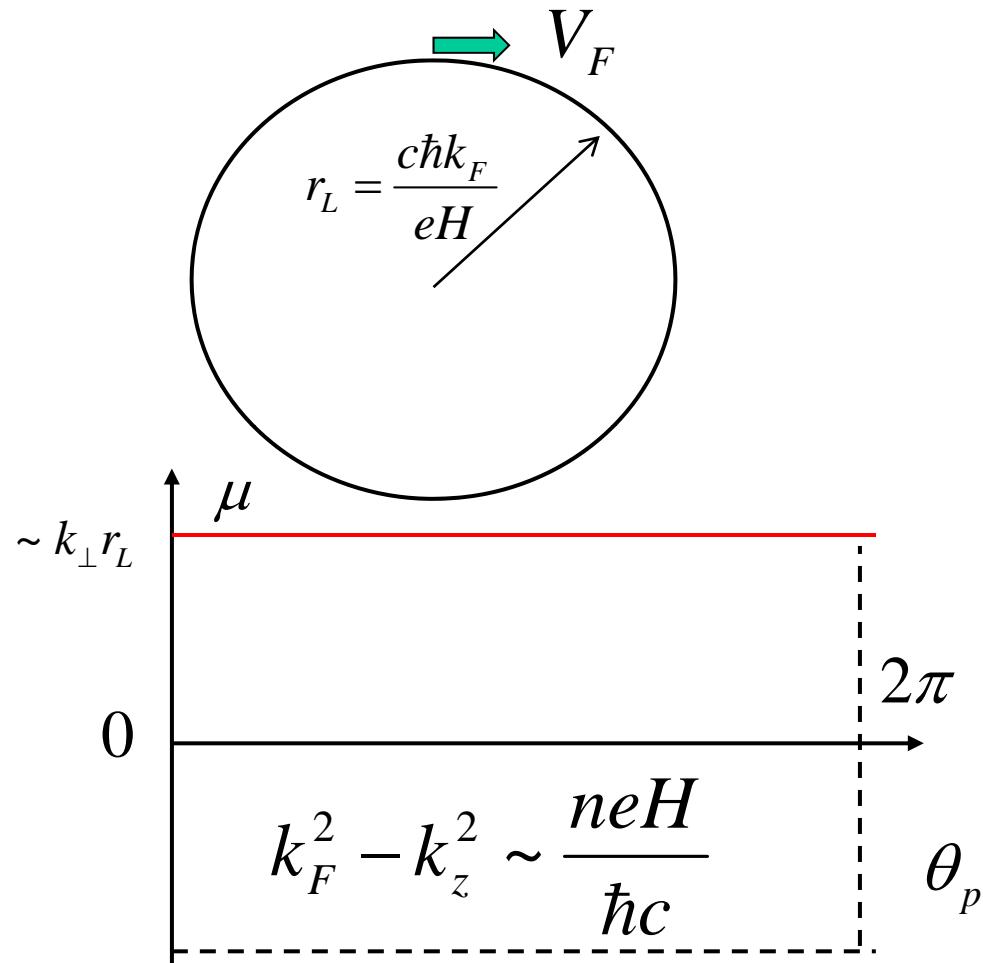


Superconductor.
Caroli – de Gennes – Matricone
spectrum with minigap.
Cyclotron orbits are destroyed.

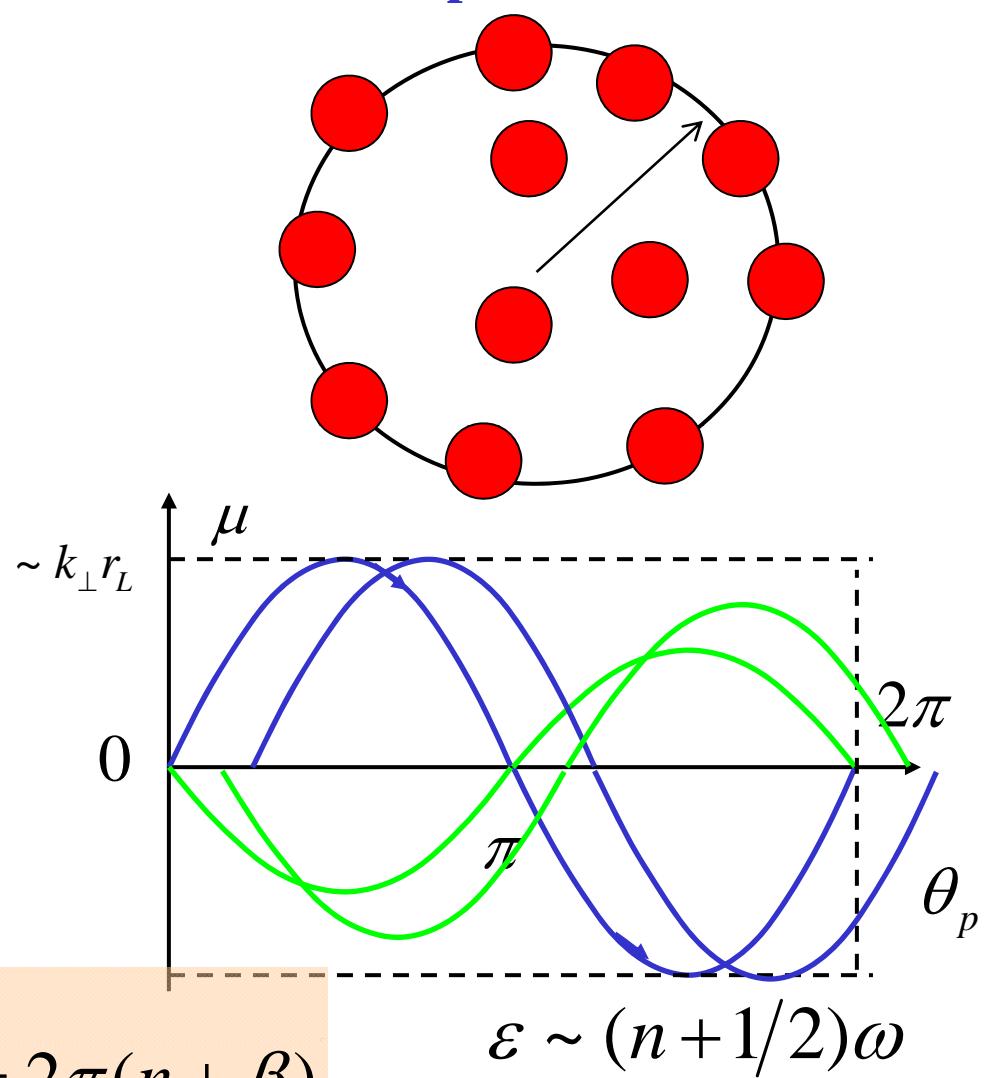


Simple example of topological transitions: de Haas-van Alfen oscillations.

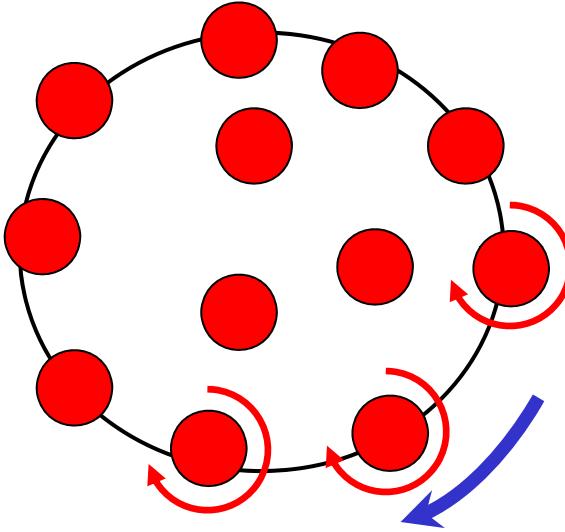
Normal metal



Superconductor.



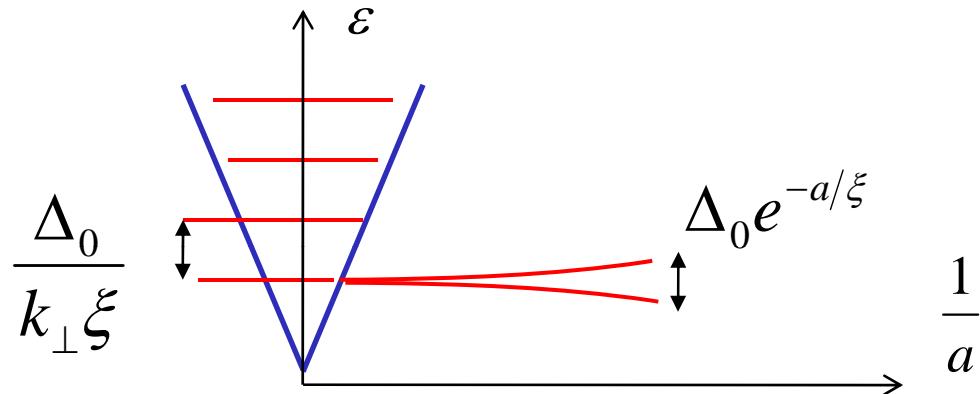
**Precession of classical trajectory inside vortex core
competes
with rotation along the cyclotron trajectory**



Question: how to restore the cyclotron motion?

Answer: to take account of quasiparticle tunneling between the vortex cores.

Qualitative arguments.
Critical intervortex distance:
minigap = energy level splitting due to tunneling



$$a_c \approx \frac{\xi}{2} \ln(k_F \xi)$$

Typical
intervortex
distance

$$k_F \xi \sim 10^2 \div 10^3$$

$$\frac{a_c}{\xi} \approx 2 \div 3$$

$$H^* \sim \frac{\phi_0}{a_c^2} \sim \frac{H_{c2}}{[\ln(k_F \xi)]^2}$$

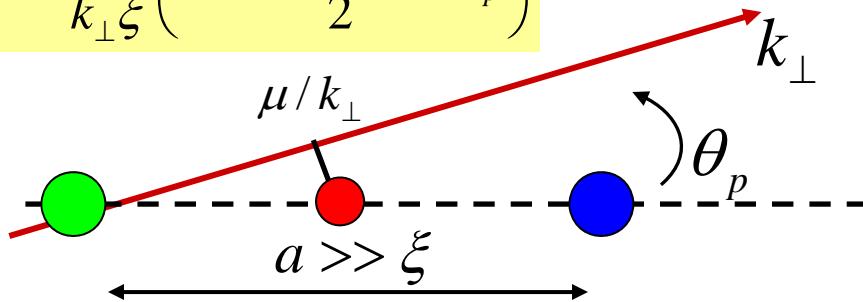
$$a \sim \sqrt{\frac{\phi_0}{H}}$$

Examples of reconnection of
Fermi lines:

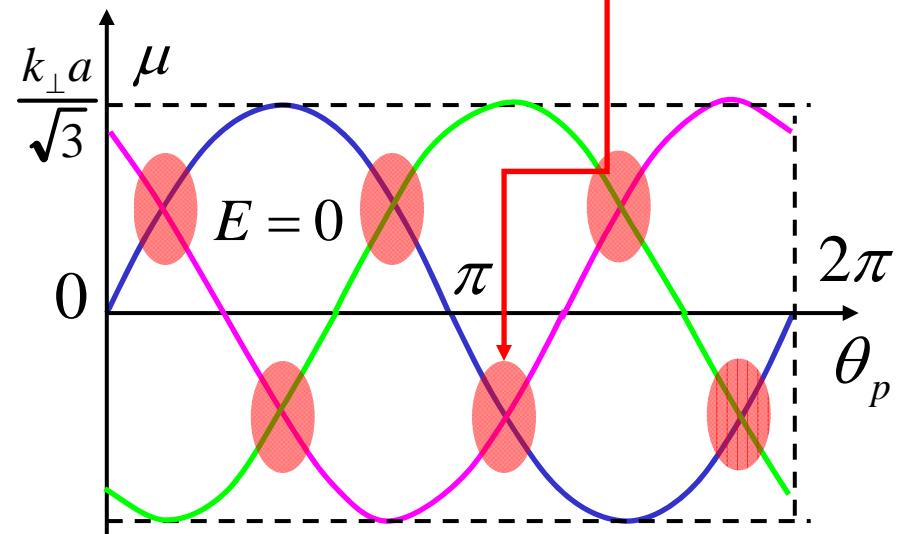
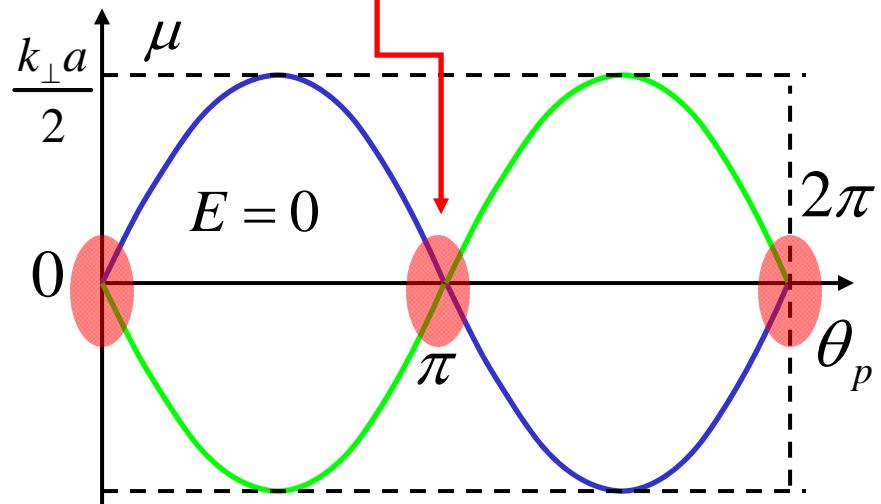
2 vortices.

3 vortices.

$$\varepsilon_{1,2} = \frac{\Delta}{k_{\perp}\xi} \left(-\mu \pm \frac{k_{\perp}a}{2} \sin \theta_p \right)$$



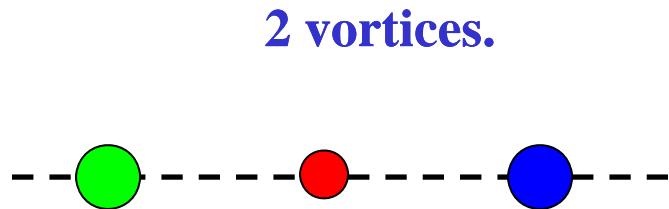
Trajectories pass through neighboring vortex cores



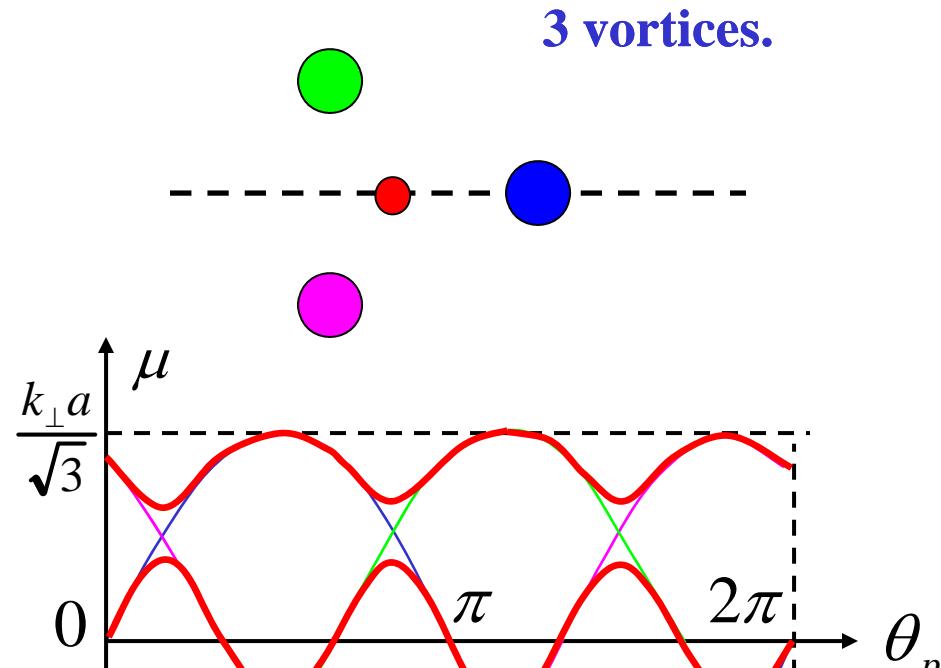
Tunneling between vortices. Splitting of CdGM energy branches.

$$\delta E \sim \Delta_0 \exp\left(-\frac{k_F a_{ij}}{k_\perp \xi}\right)$$

$$a_{ij} = |\vec{r}_i - \vec{r}_j|$$

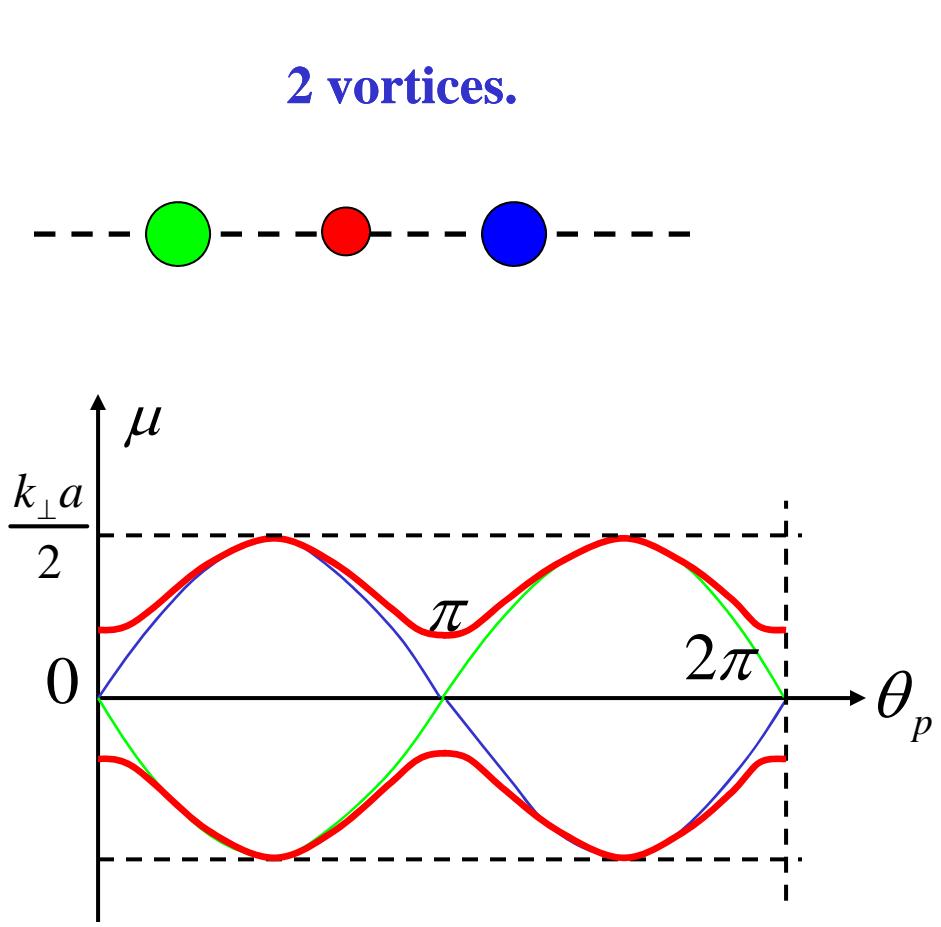


$$\delta\mu(a_{ij}) \sim \frac{\delta E}{\omega} \sim k_\perp \xi \exp\left(-\frac{k_F a_{ij}}{k_\perp \xi}\right)$$

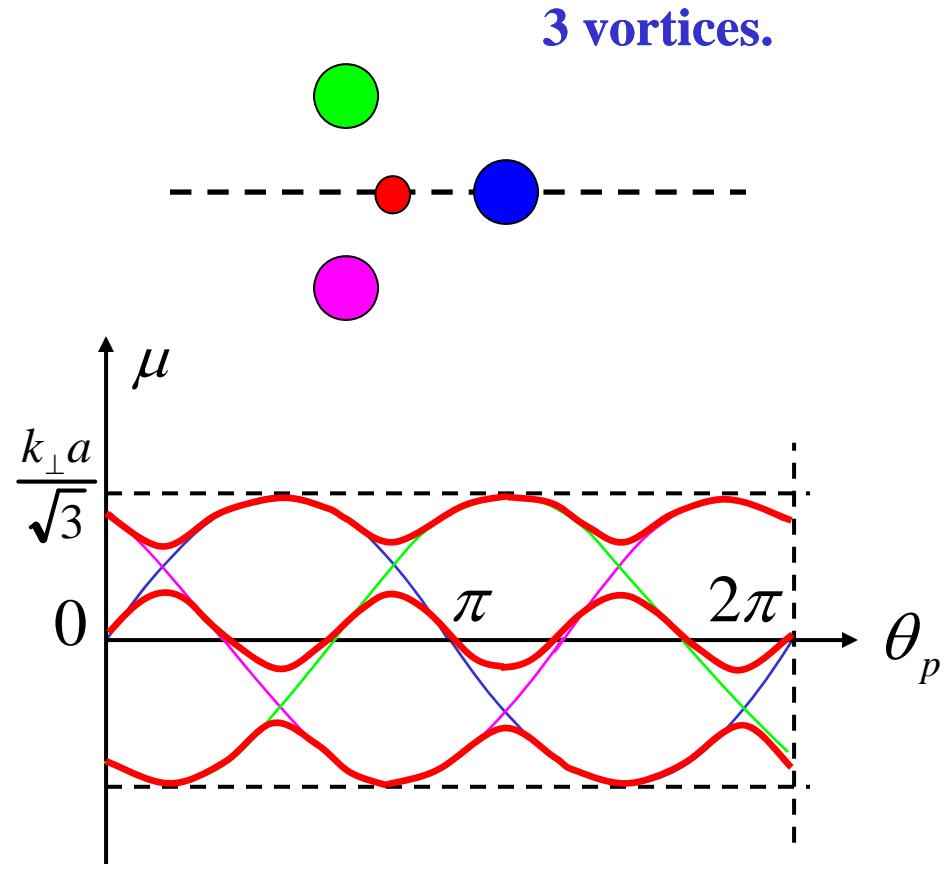


Crossover to a multiquantum vortex

2 vortices.

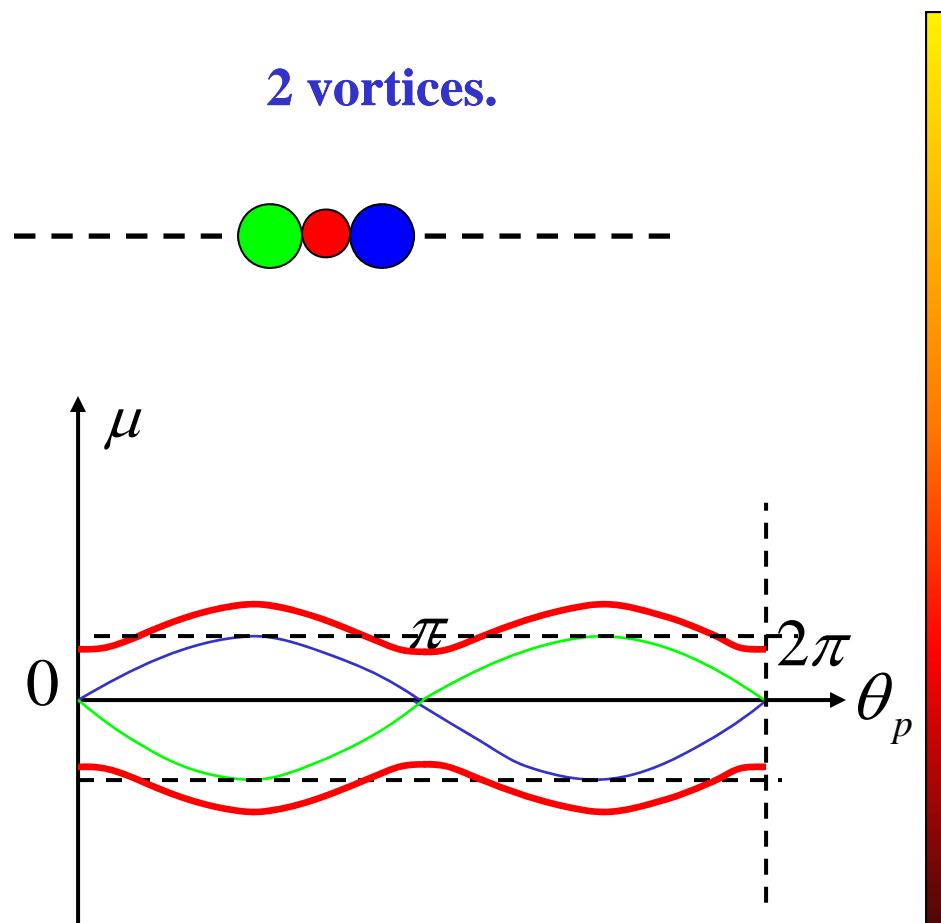


3 vortices.

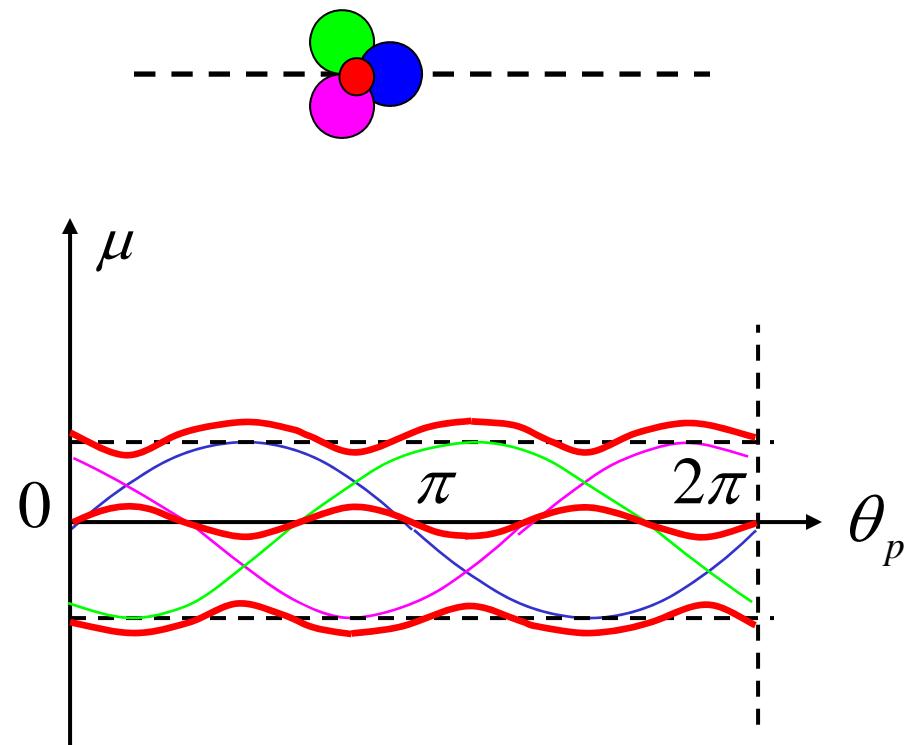


Crossover to a multiquantum vortex

2 vortices.



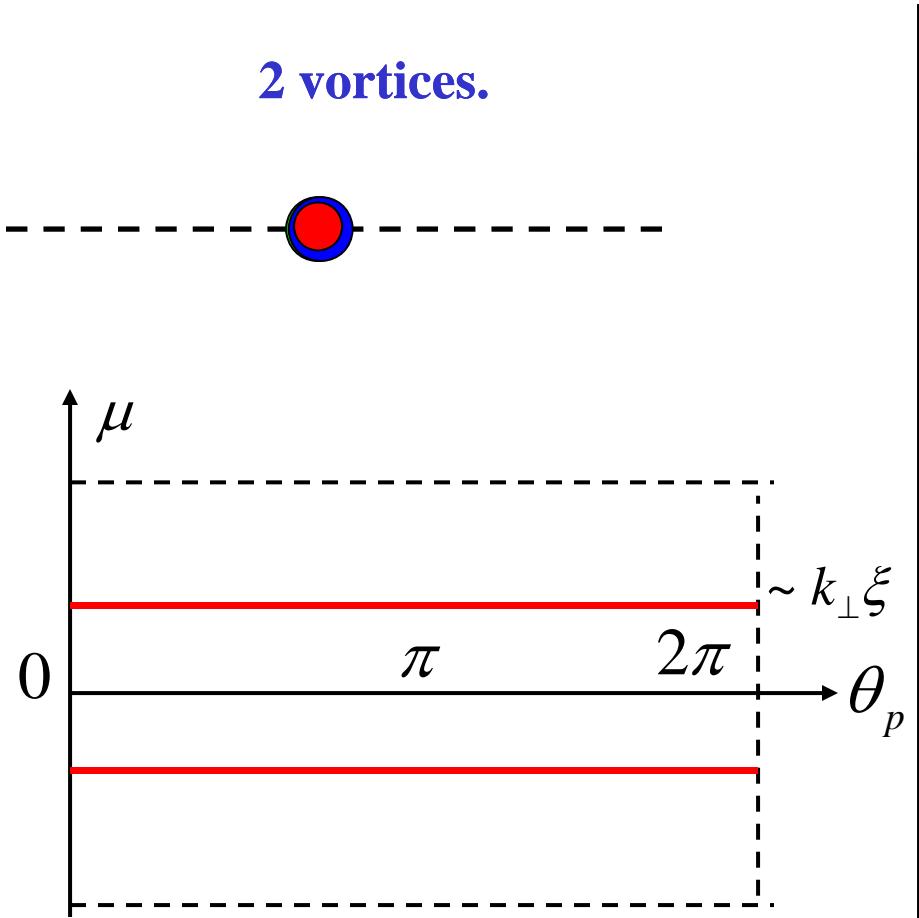
3 vortices.



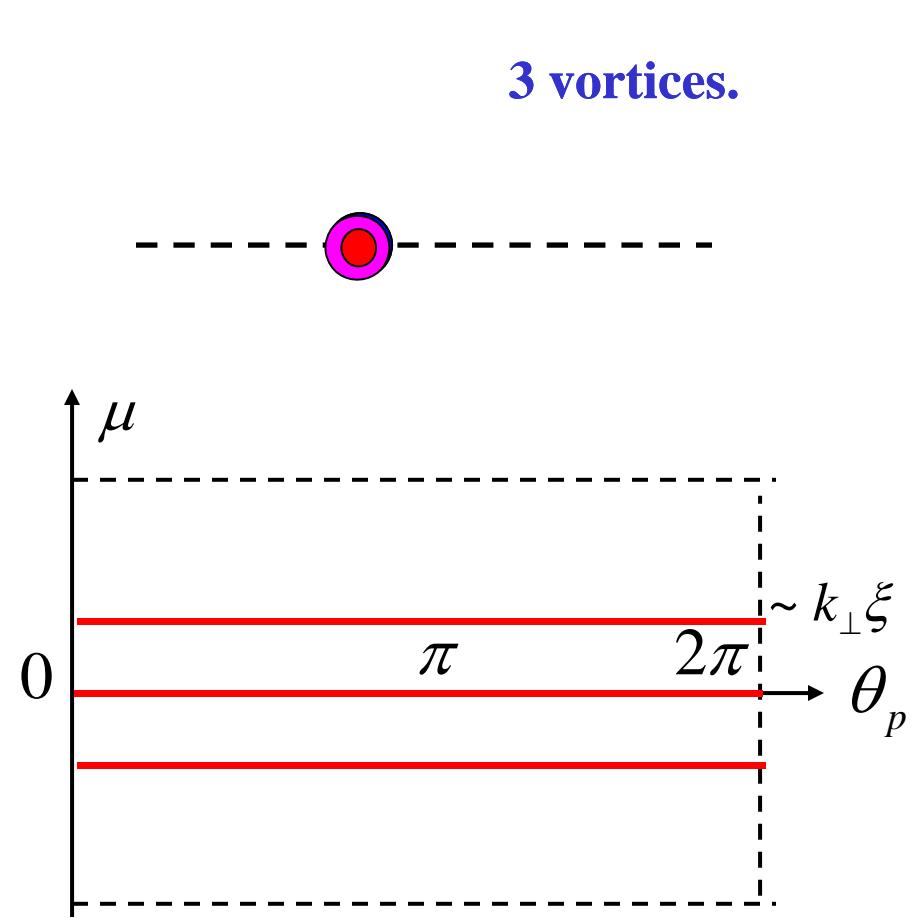
Crossover to a multiquantum vortex

$$a = 0$$

2 vortices.



3 vortices.



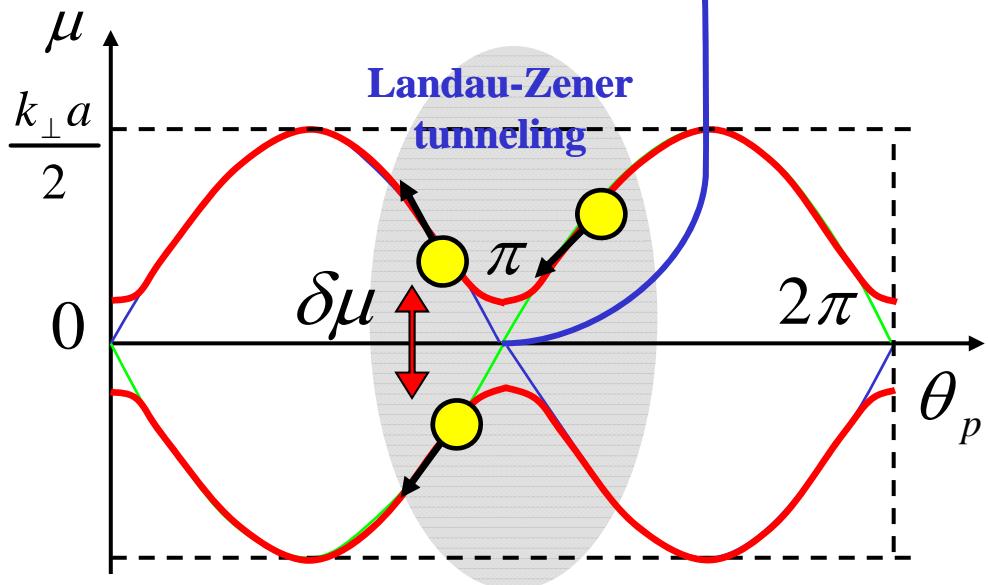
Quantum mechanics of precessing trajectories

$$[\theta_p, \hat{\mu}] = i$$

Uncertainty principle: $\Delta\mu\Delta\theta_p \sim 1$

$$\Delta\mu \sim k_{\perp} a_{ij} \Delta\theta_p$$

Quantum mechanical uncertainty of the angular momentum near the branch crossing points – $\Delta\mu \sim \sqrt{k_{\perp} a_{ij}}$



Criterion of intervortex tunneling efficiency:

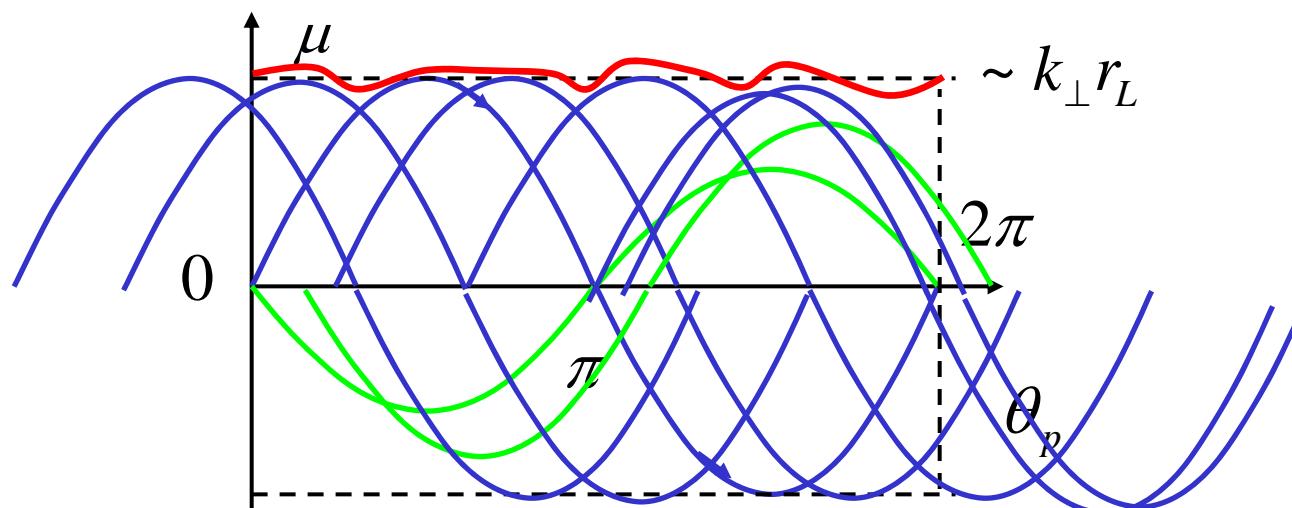
$$\delta\mu \gg \Delta\mu$$

$$\frac{k_{\perp}\xi^2}{a_{ij}} \exp\left(-\frac{2k_F a_{ij}}{k_{\perp}\xi}\right) \gg 1$$

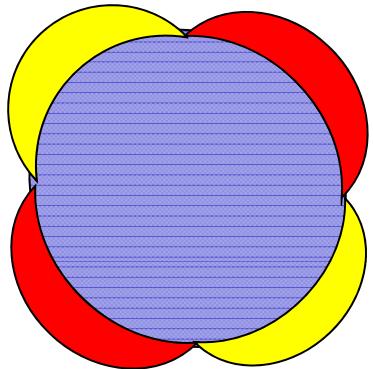
Reconnection of Fermi surface segments due to the intervortex tunneling

Magnetic breakdown criterion.

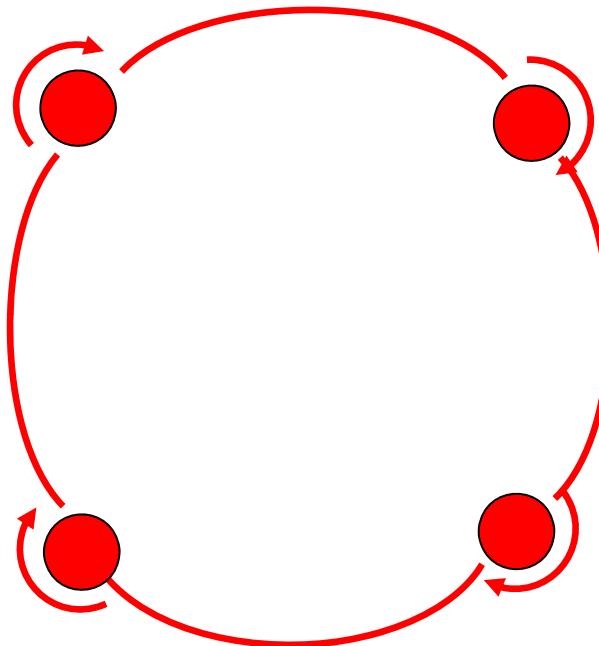
$$H > H^* \sim \frac{H_{c2}}{[\ln(k_F \xi)]^2}$$



Useful generalization: d-wave pairing.
Cyclotron arcs are connected by
precession of trajectories inside the cores



$$\Delta = \frac{\Delta_0 k_x k_y}{k_F^2}$$

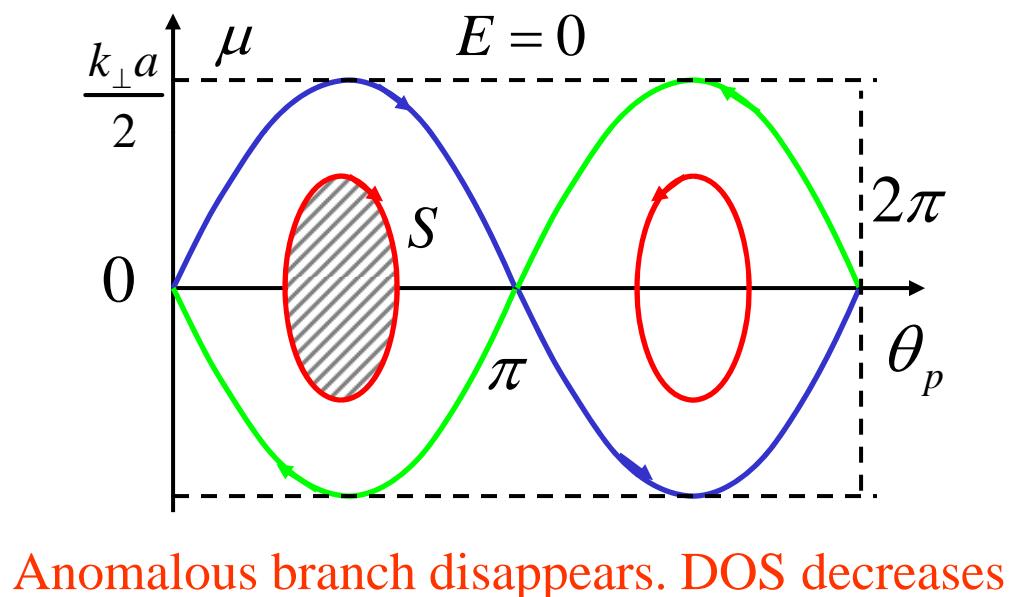
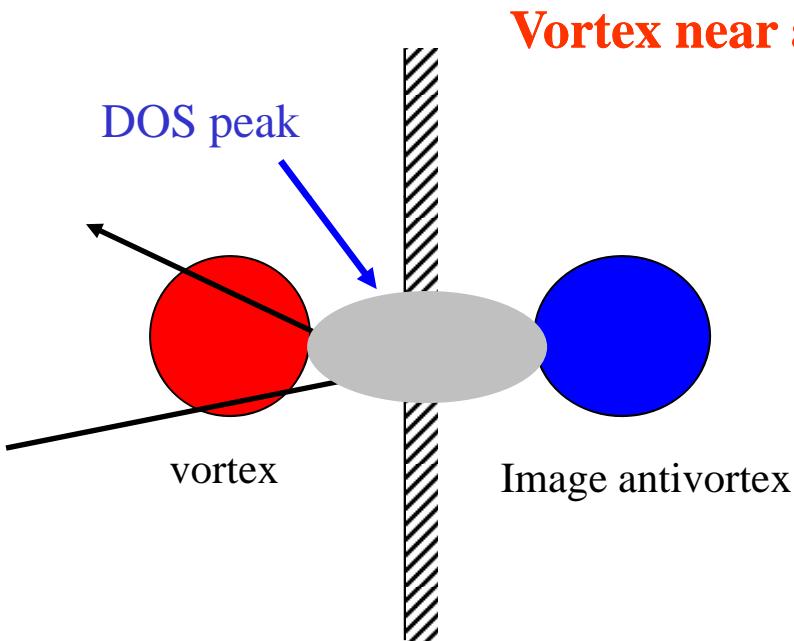


$$\delta\theta \sim \frac{a_c}{a}$$

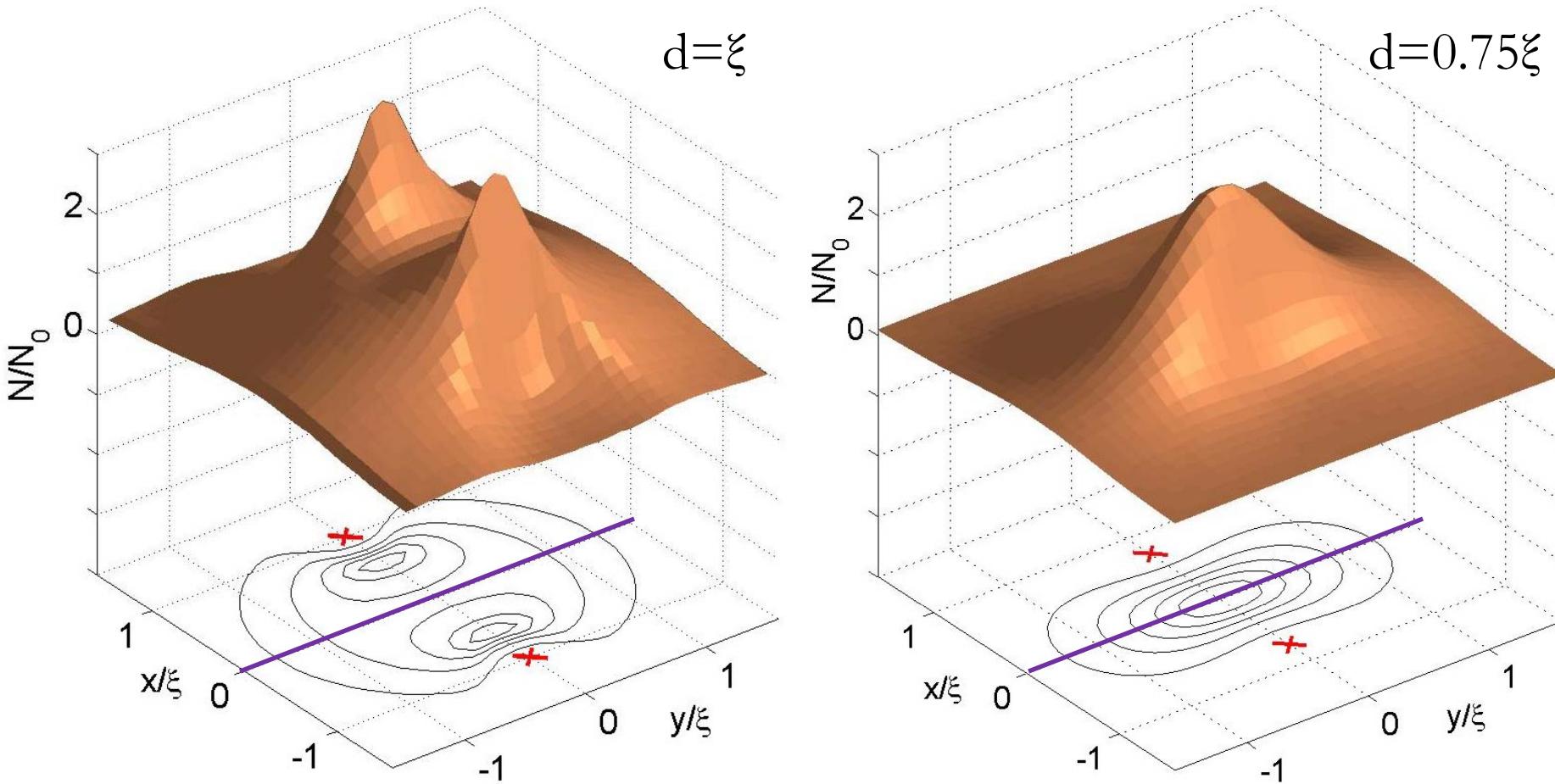
$$\int_0^{2\pi n_\theta} \mu(\theta_p) d\theta_p \sim k_F r_L \frac{a_c}{a} \sim 2\pi(n + \beta)$$

$$k_F^2 \sqrt{\frac{H}{H^*}} \sim \frac{neH}{\hbar c}$$

Examples of topological transitions.
Opening of Fermi surface segments in vortex cores.



Local DOS for a vortex positioned near the surface



LDOS peak is shifted towards the boundary

Several boundaries: splitting of LDOS peak becomes possible

Vortex pinning by columnar defects

Artificial columnar pinning centers:

Proton or heavy-ion irradiation : $2R \sim 50\text{\AA}$

Inclusion of normal particles or nanorods: $2R \sim 50\text{-}100\text{\AA}$

Submicrometer holes (nanorods): $2R \sim 0.1\text{-}1 \mu\text{m}$

Attraction of vortex to a cavity (hole)

G.S.Mkrtchyan and V.V.Shmidt (1972)

A.Buzdin and D.Feinberg (1996)

A.Buzdin and M.Daumens (1998, 2000)

H.Nordborg and V.Vinokur (2000)

Microscopic theory of vortex pinning by scattering centers

E.V.Thuneberg et al. (1982,1984)

$$\sigma \ll \xi_0^2$$

Multiquanta vortices trapped by the cavities

A.I.Buzdin (1993)

A.Bezryadin, A.I. Buzdin, B. Pannetier (1994)

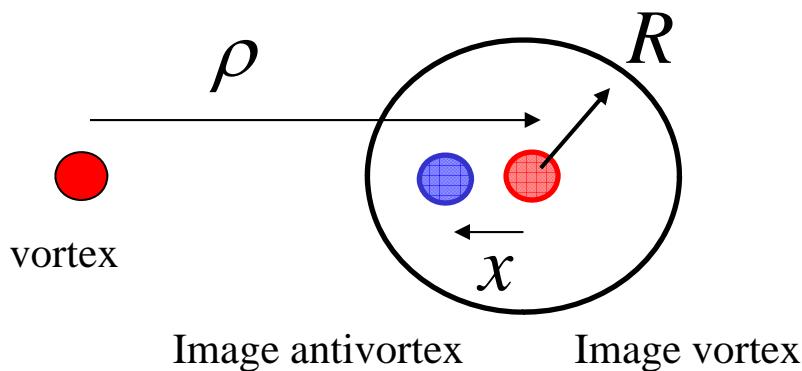
A.Bezryadin, B. Pannetier (1995)

V.Bruyndoncx et al (1999)

Pinning mechanisms

$$R \gg \xi$$

Electrodynamic mechanism



$$x\rho = R^2$$

$$R \ll \xi$$

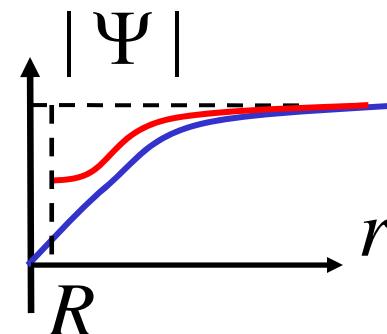
Gain in the vortex core energy

Local contribution

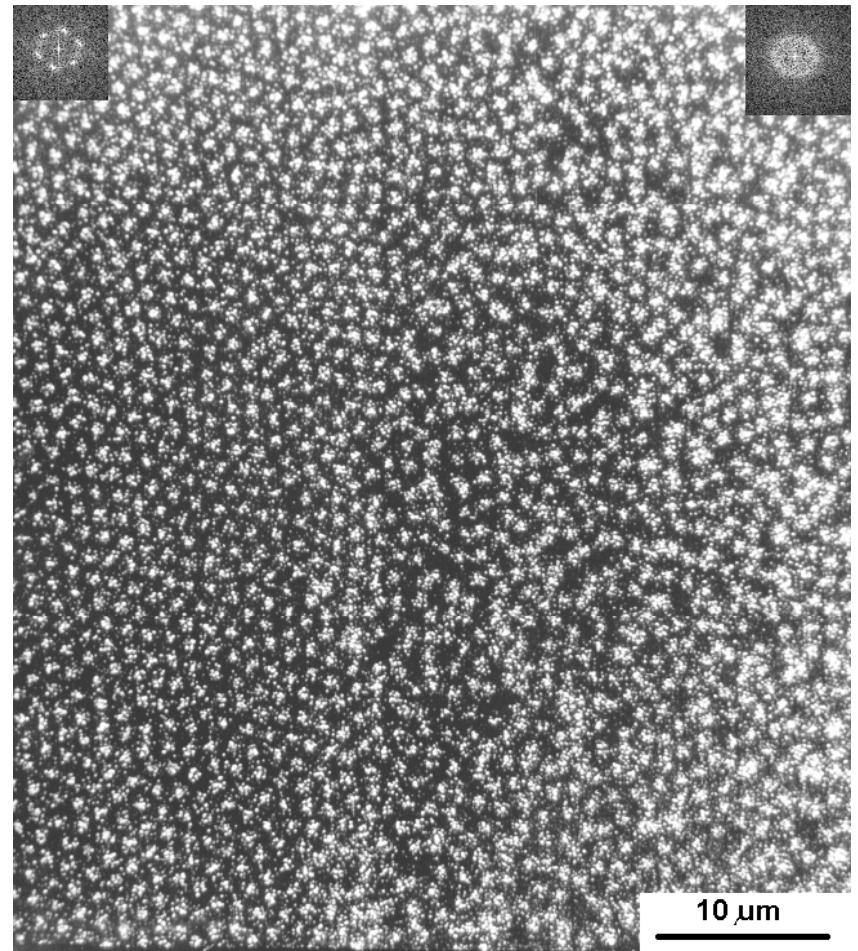
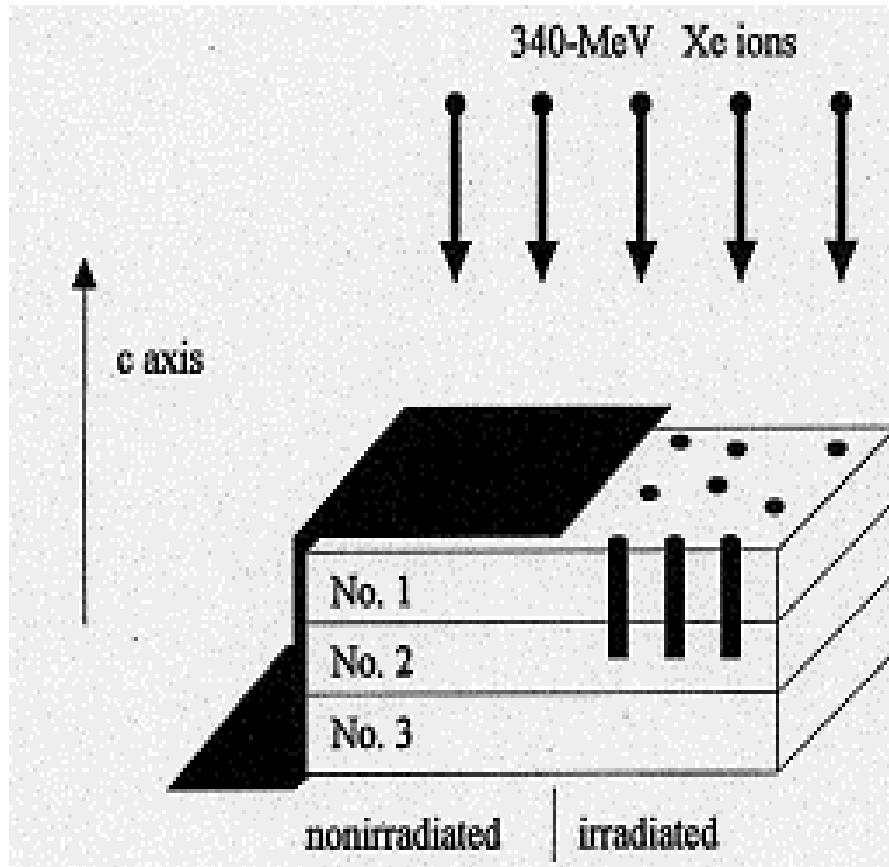
$$\delta E \sim \frac{H_{cm}^2}{8\pi} V_{defect}$$

Nonlocal contribution

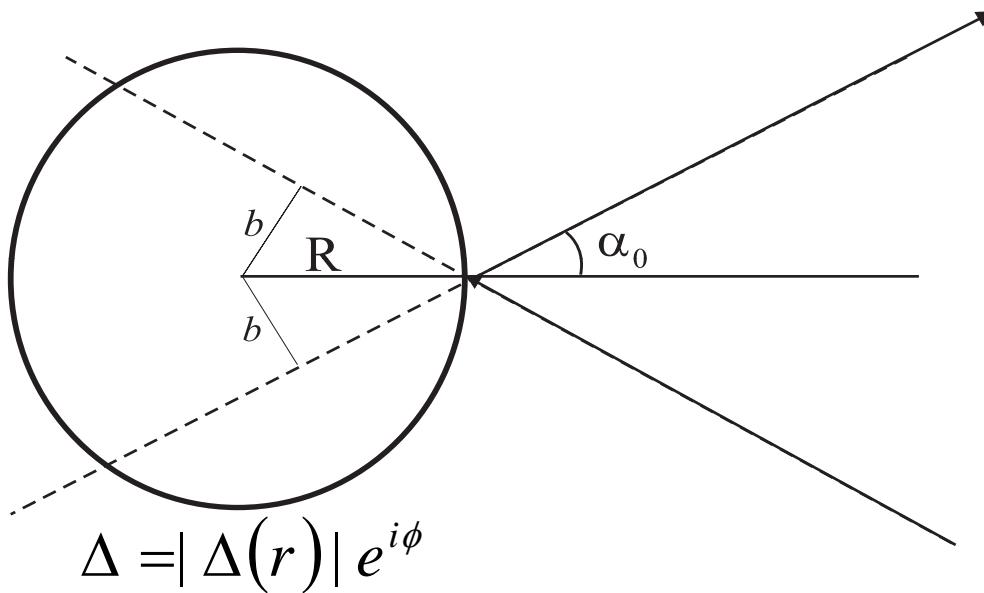
$$\delta E > \frac{H_{cm}^2}{8\pi} V_{defect}$$



Flux Line Lattice destruction by columnar defects in BSCCO



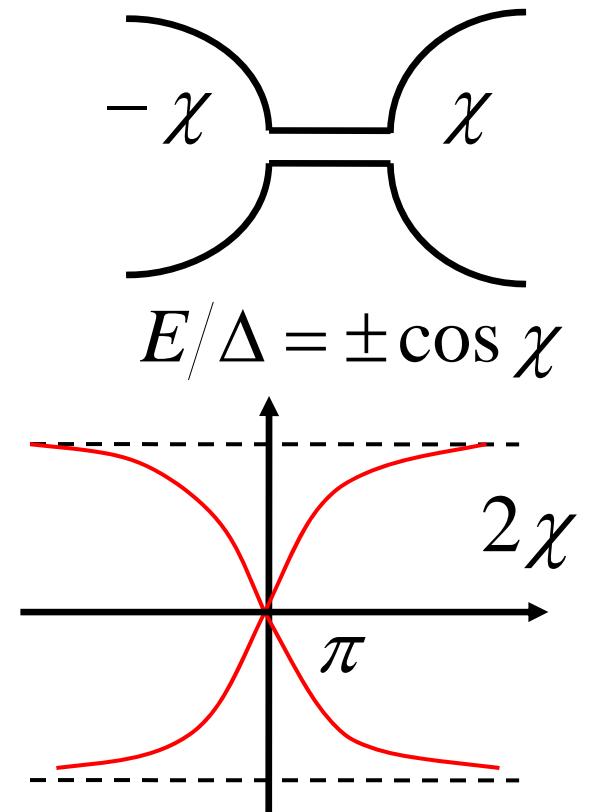
Vortex trapped by a defect with radius $R \gg \lambda_F$
Analogy with quasiparticle spectrum of Josephson junction



Boundary condition: $\hat{\Psi}(R) = 0$

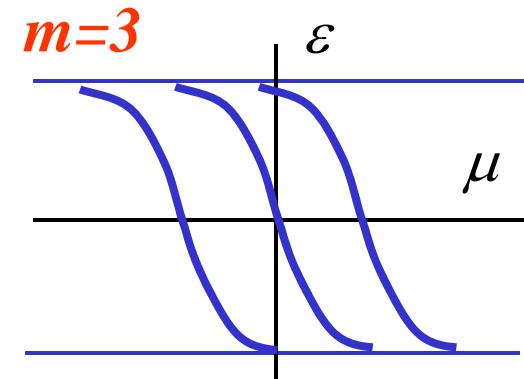
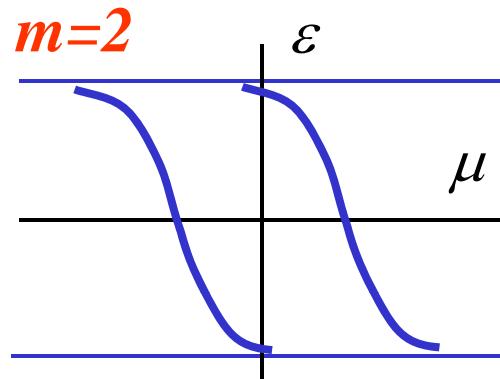
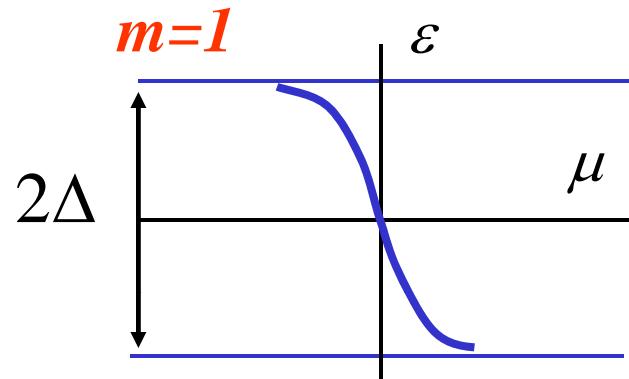
$$E/\Delta_0 = \cos \chi = \cos \alpha_0$$

$$E/\Delta_0 = \text{sign}(b) \sqrt{1 - \frac{b^2}{R^2}}$$

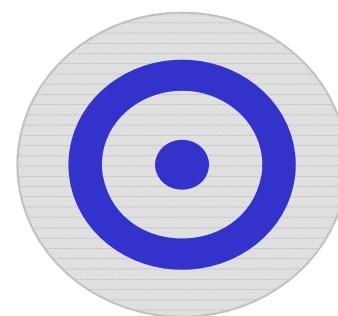
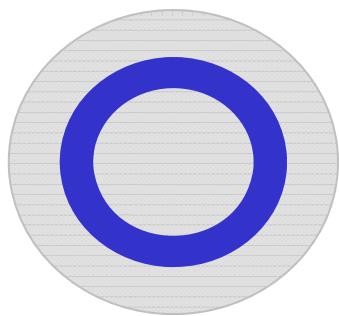
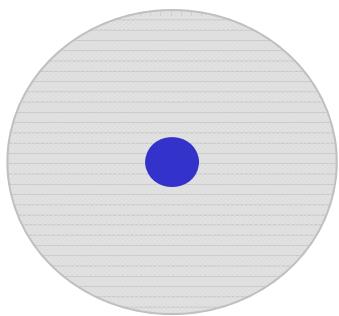


Anomalous spectral branches for different vorticities.

spectra

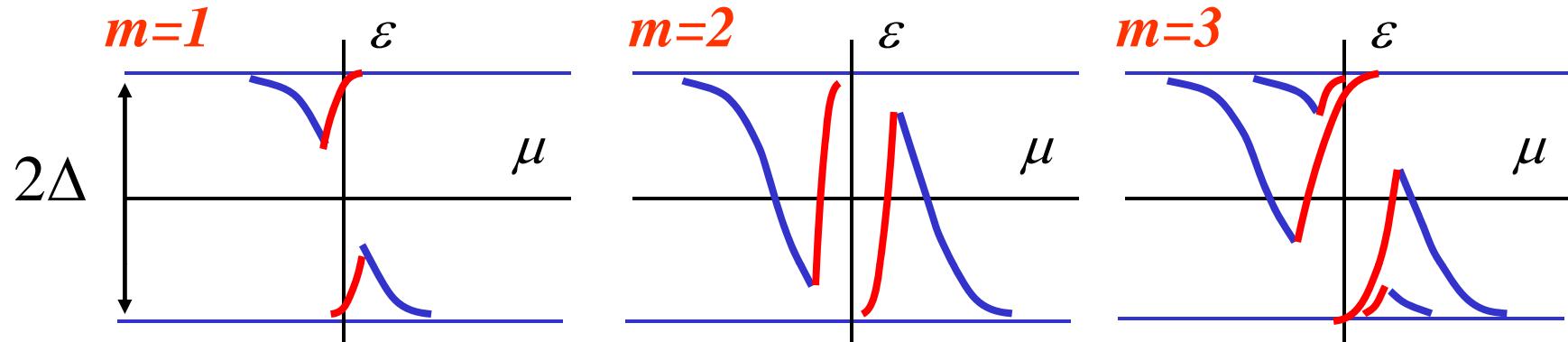


DOS

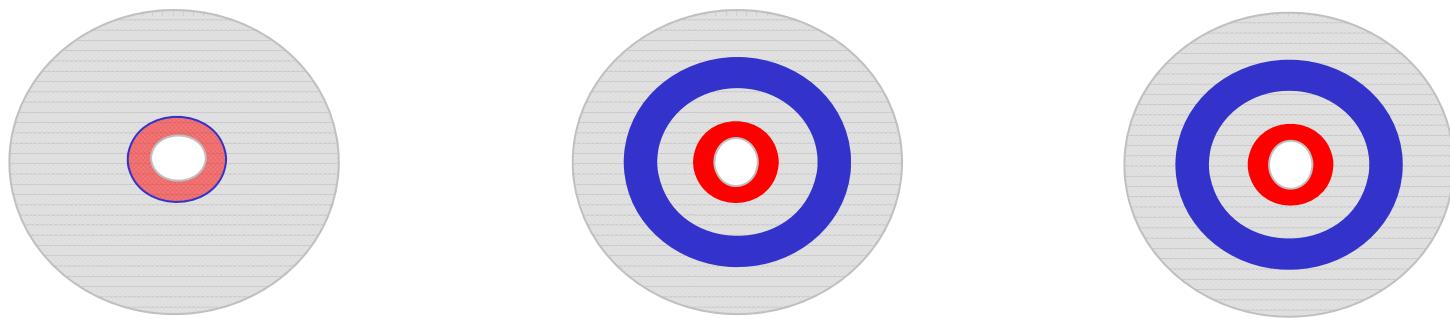


Anomalous spectral branches for different vorticities. Influence of defect

spectra



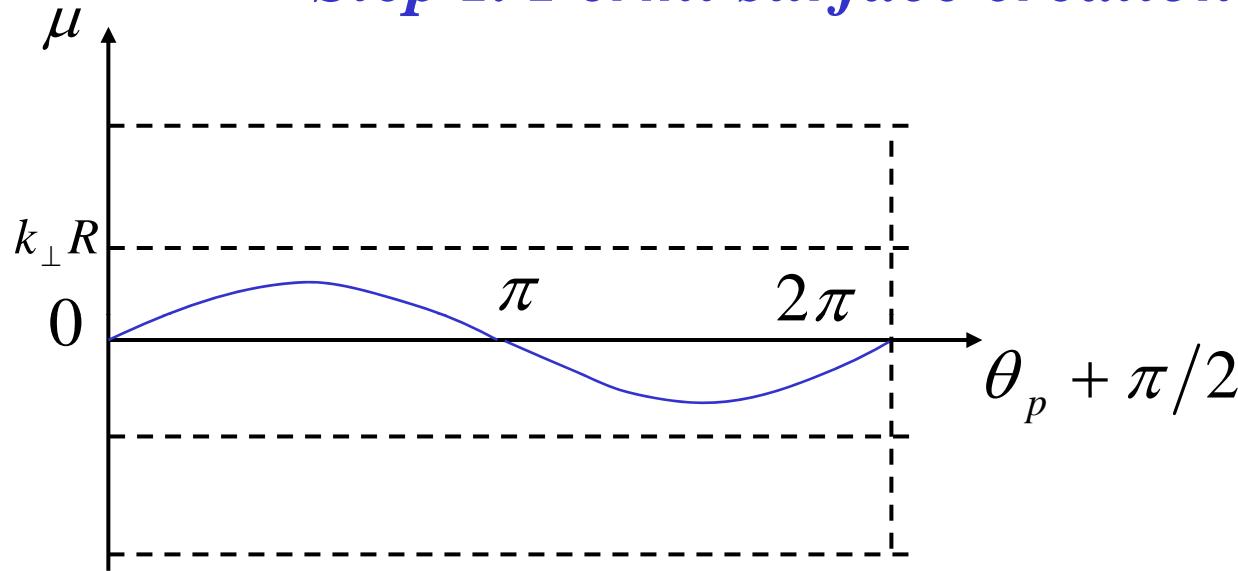
DOS



Increasing minigap

Vortex escape as a Fermi surface transformation

Step I. Fermi surface creation

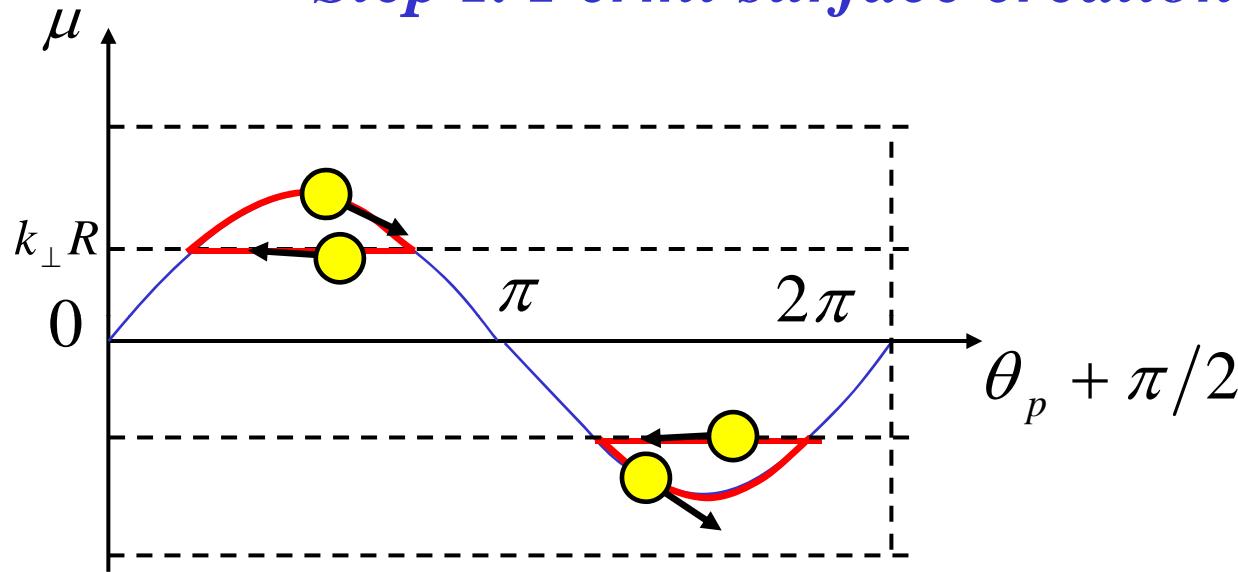


$$E = -\omega\mu + \hbar k_{\perp} V_s \cos \theta_p$$

$$\mu = \frac{\hbar k_{\perp} V_s}{\omega} \cos \theta_p \qquad V_s < \frac{\omega R}{\hbar}$$

Vortex escape as a Fermi surface transformation

Step I. Fermi surface creation



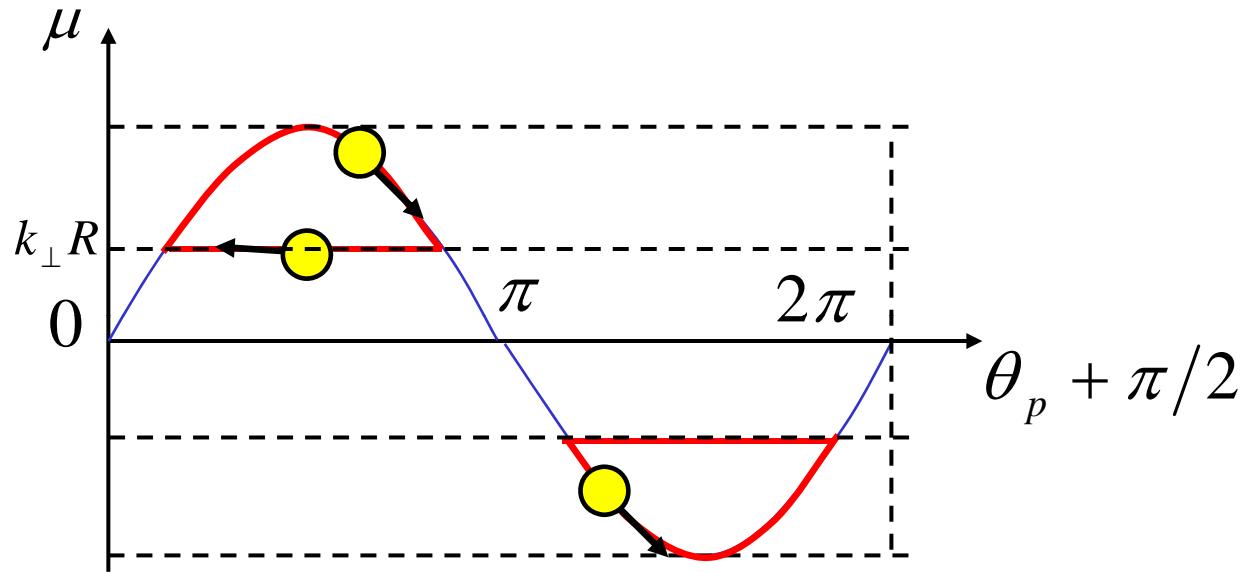
$$E = -\omega\mu + \hbar k_{\perp} V_s \cos \theta_p$$

$$\mu = \frac{\hbar k_{\perp} V_s}{\omega} \cos \theta_p$$

$$V_s > \frac{\omega R}{\hbar}$$

Vortex escape as a Fermi surface transformation

Step I. Fermi surface creation. Vortex is bound to the cavity



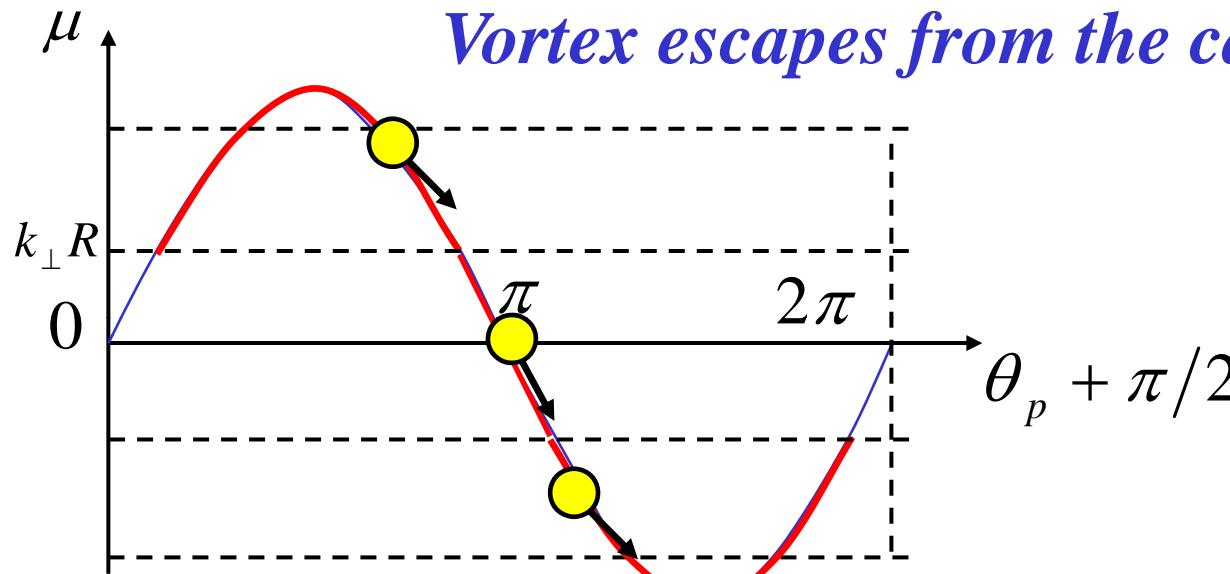
$$E = -\omega\mu + \hbar k_{\perp} V_s \cos \theta_p$$

$$\mu = \frac{\hbar k_{\perp} V_s}{\omega} \cos \theta_p$$

$$V_s > \frac{\omega R}{\hbar}$$

Vortex escape as a Fermi surface transformation

*Step II. Merging of two parts of Fermi surface.
Vortex escapes from the cavity.*



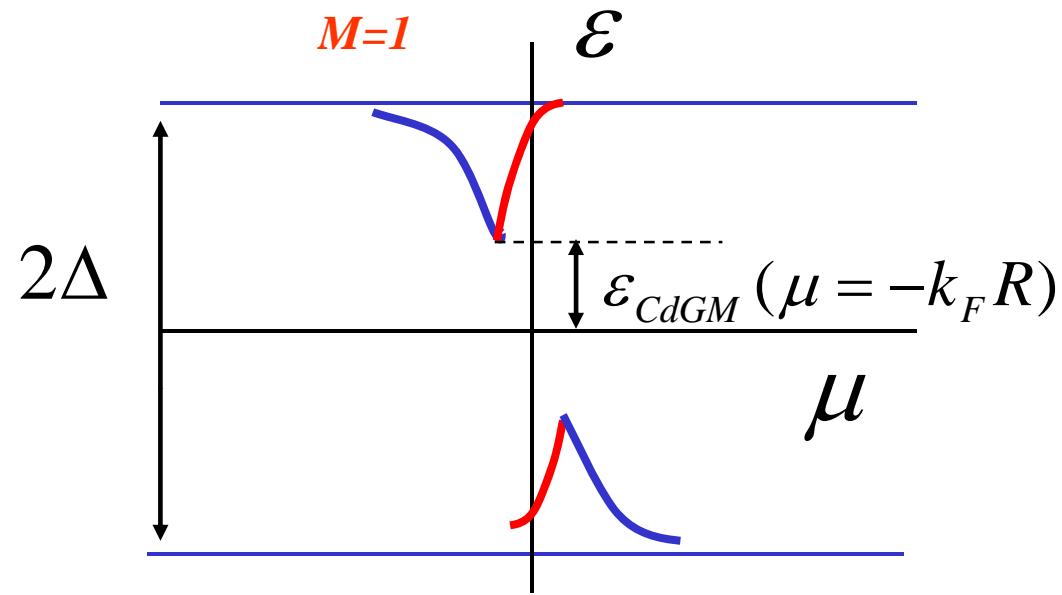
$$E = -\omega\mu + \hbar k_{\perp} V_s \cos \theta_p$$

$$\mu = \frac{\hbar k_{\perp} V_s}{\omega} \cos \theta_p$$

$$V_s > \frac{\omega R}{\hbar}$$

$$\Delta\mu > k_{\perp} R$$

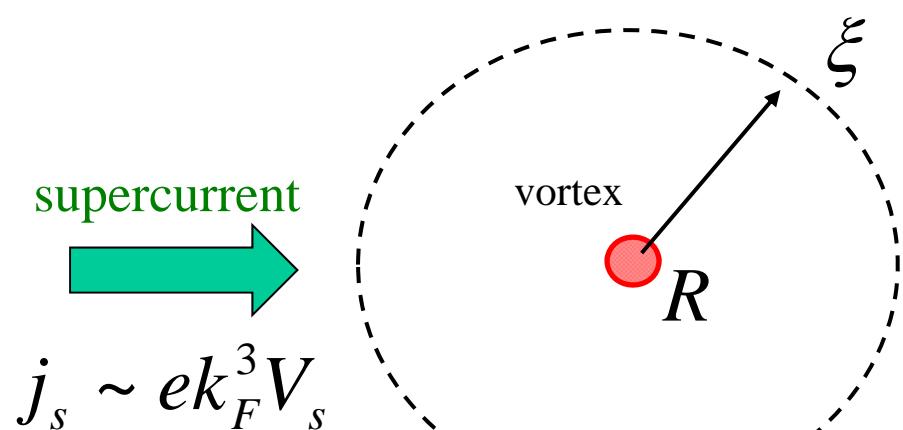
Depinning current I.



Small defect radius

$$R \ll \xi$$

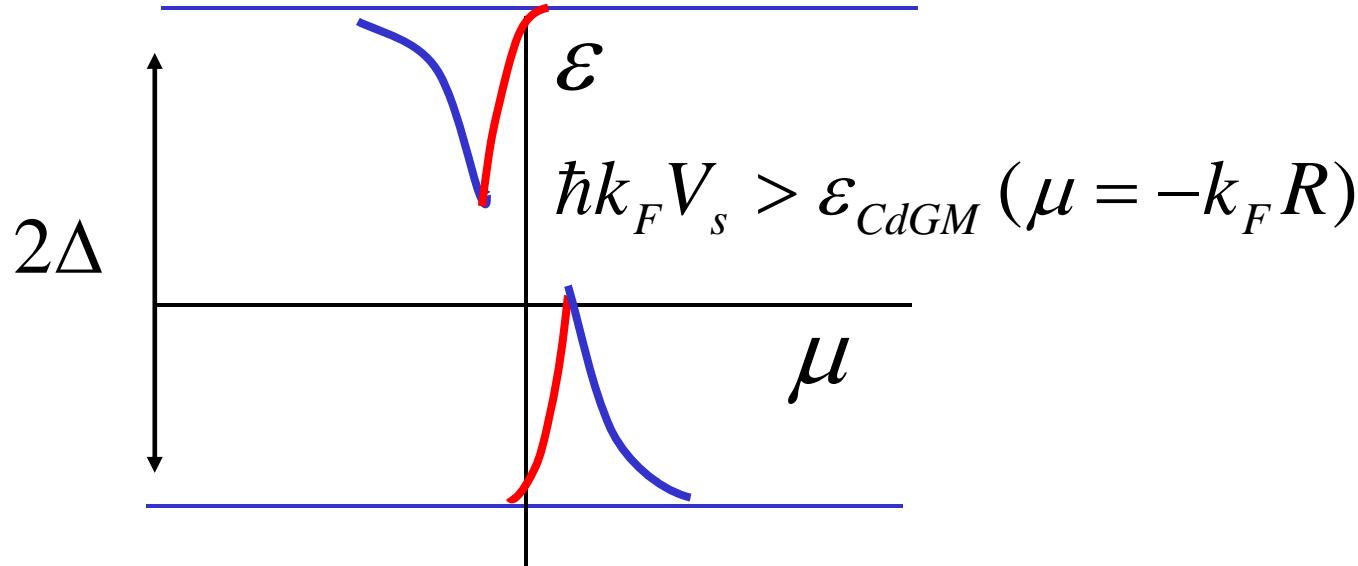
$$\epsilon_{CdGM}(\mu = -k_F R) \approx \Delta \frac{R}{\xi}$$



Doppler shift of quasiparticle energy:

$$\epsilon = \epsilon_{CdGM}(\mu = -k_F R) + \hbar \vec{k}_F \vec{V}_s$$

Depinning current I.



Critical superfluid velocity: $V_{c1} = \varepsilon_{CdGM} (\mu = -k_F R) / \hbar k_F$

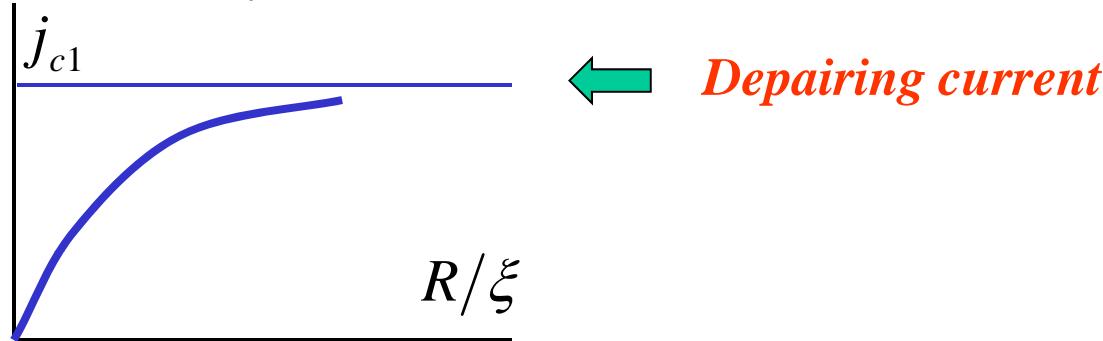
Critical supercurrent: $j_{c1} = \frac{e}{\hbar} k_F^2 \varepsilon_{CdGM} (\mu = -k_F R)$

Small defect radius $j_{c1} \approx \frac{e}{\hbar} k_F^2 \Delta \frac{R}{\xi}$

Large defect radius $j_{c1} \approx \frac{e}{\hbar} k_F^2 \Delta$ - Depairing current

Depinning current I.

$$j_{c1} = \frac{e}{\hbar} k_F^2 \mathcal{E}_{CdGM} (\mu = -k_F R)$$



Comparison with textbook depinning current estimates for $R < \xi$

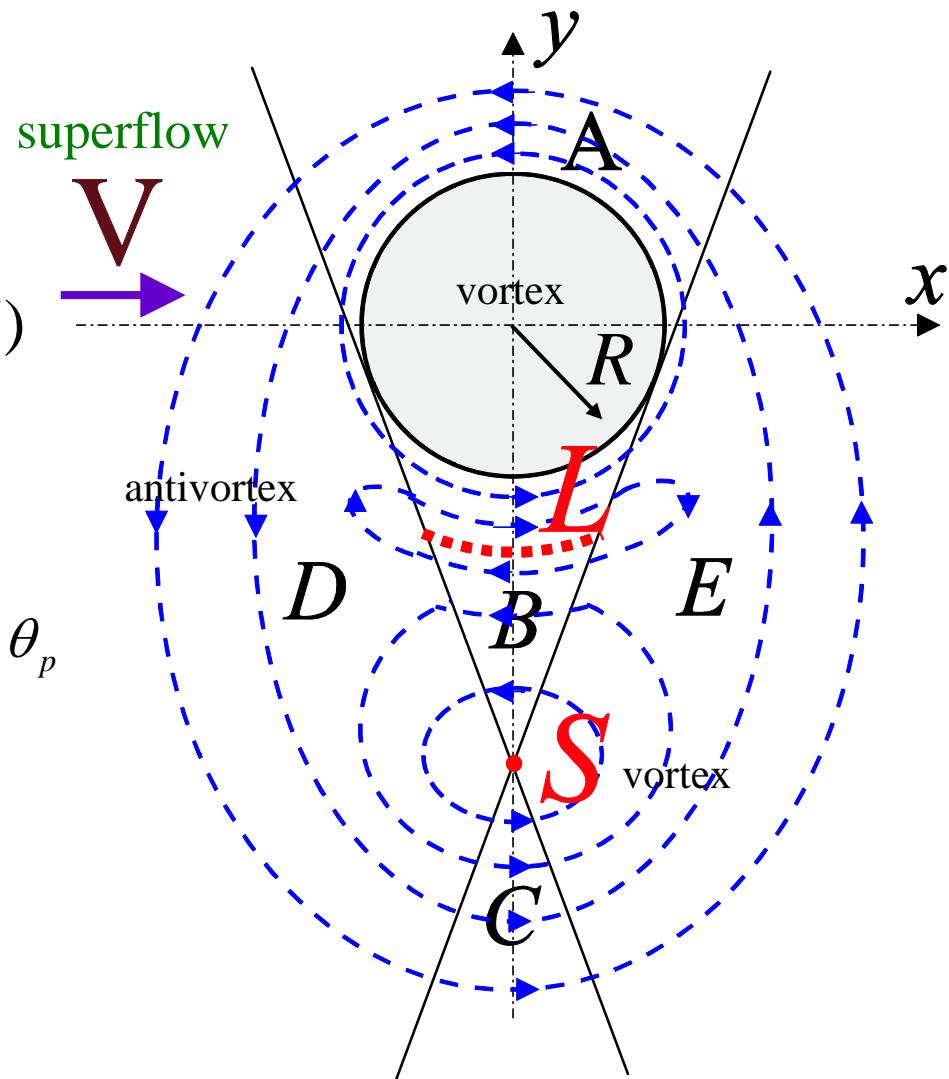
Energy gain for a vortex pinned at the cylindrical defect: $\delta U \sim H_{cm}^2 R^2$

Depinning current: $j_p = \frac{\delta U}{\xi} \sim H_{cm}^2 R \frac{R}{\xi} \sim j_{c1} \frac{R}{\xi} \ll j_{c1}$

Superconducting gap profile. Scenario of the vortex escape from a small cavity.

$$\frac{\Delta_{core}(r, \theta)}{\Delta_\infty} = \int_0^{2\pi} \frac{d\theta_p}{2\pi} e^{i\theta_p} \tanh\left(\frac{E}{2T}\right) h(E)$$

$$E(r, \theta, \theta_p) = \varepsilon(b = r \sin(\theta_p - \theta)) + \hbar k_F V_s \cos \theta_p$$

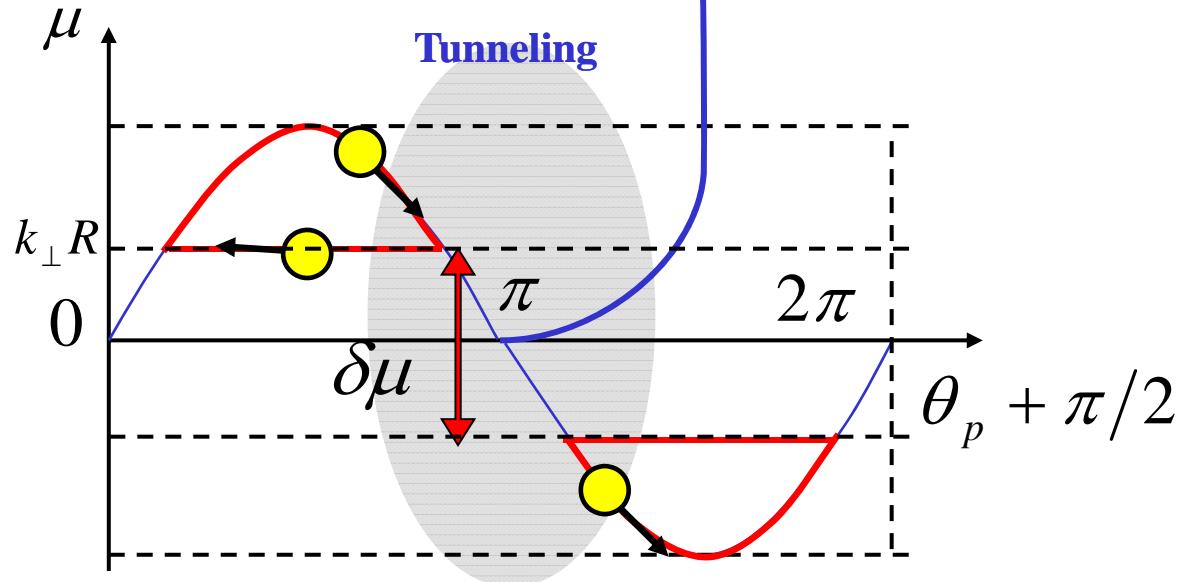


Depinning current II.

$$[\theta_p, \hat{\mu}] = i$$

Uncertainty principle: $\Delta\mu\Delta\theta_p \sim 1$

$$\Delta\mu \sim \hbar k_{\perp} V_s \Delta\theta_p / \omega$$



Quantum mechanics of precessing trajectories

Quantum mechanical uncertainty of angular momentum near the branch crossing point:

$$\Delta\mu \sim \sqrt{\frac{\hbar k_{\perp} V_s}{\omega}}$$

Criterion of tunneling efficiency:

$$\Delta\mu > k_{\perp} R$$

$$V_s > \frac{k_{\perp} R^2 \omega}{\hbar}$$

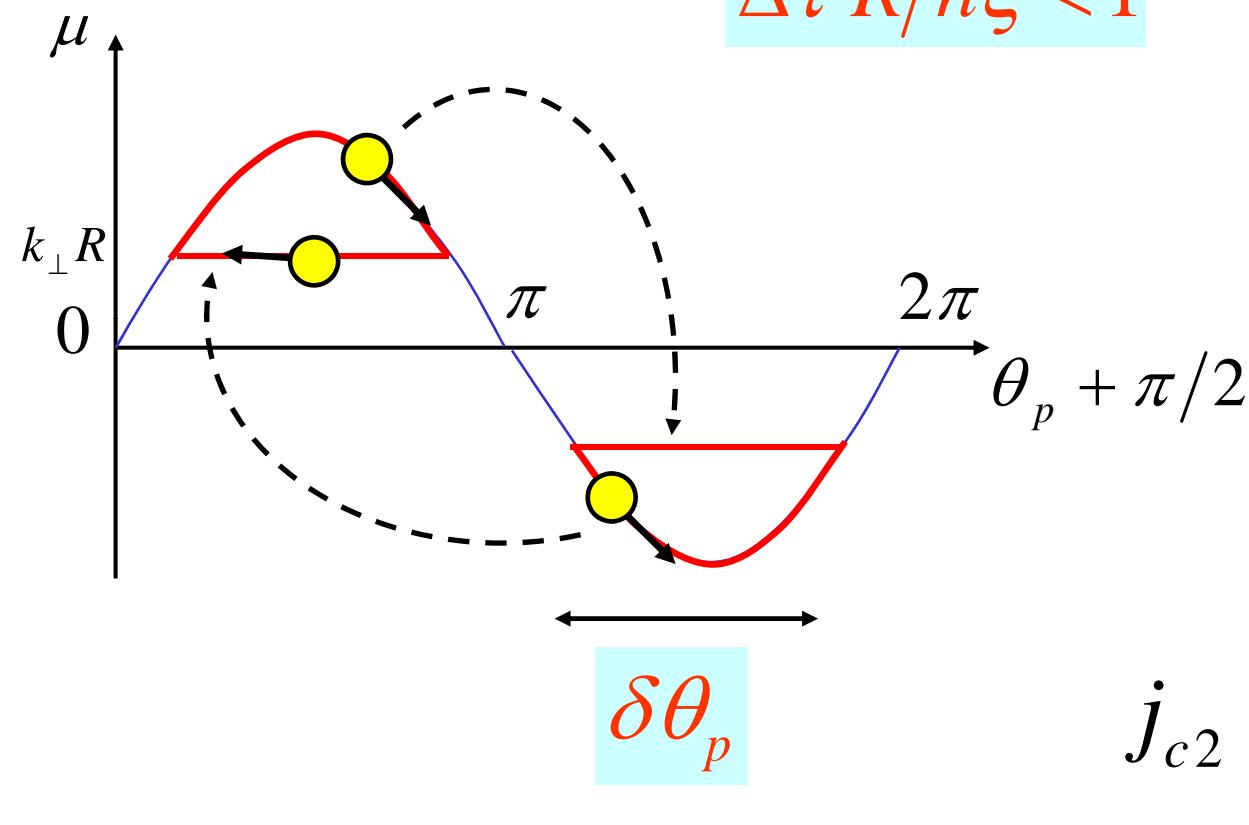
Depinning current II.

$$V_{s2} > \frac{k_F R^2 \omega}{\hbar}$$

$$j_{c2} \approx \frac{e}{\hbar} k_F^3 \Delta \frac{R^2}{\xi} \sim j_{depair} \frac{k_F R^2}{\xi} \sim j_{c1} k_F R$$

Depinning current II. Effect of impurities

$$\Delta \tau R / \hbar \xi < 1$$



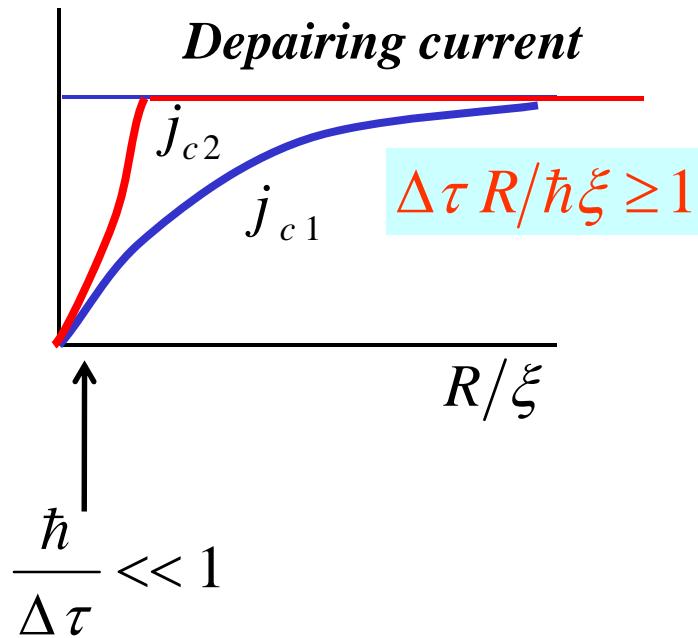
$$\frac{1}{\tau_s} \sim \frac{\delta\theta_p}{\tau} < \frac{\Delta R}{\hbar \xi}$$

$$\cos(\delta\theta_p/2) = \frac{j_{c1}}{j}$$

$$j_{c2} \approx \frac{\dot{j}_{c1}}{\cos(\gamma \Delta \tau R / \hbar \xi)}$$

Depinning current at low temperatures

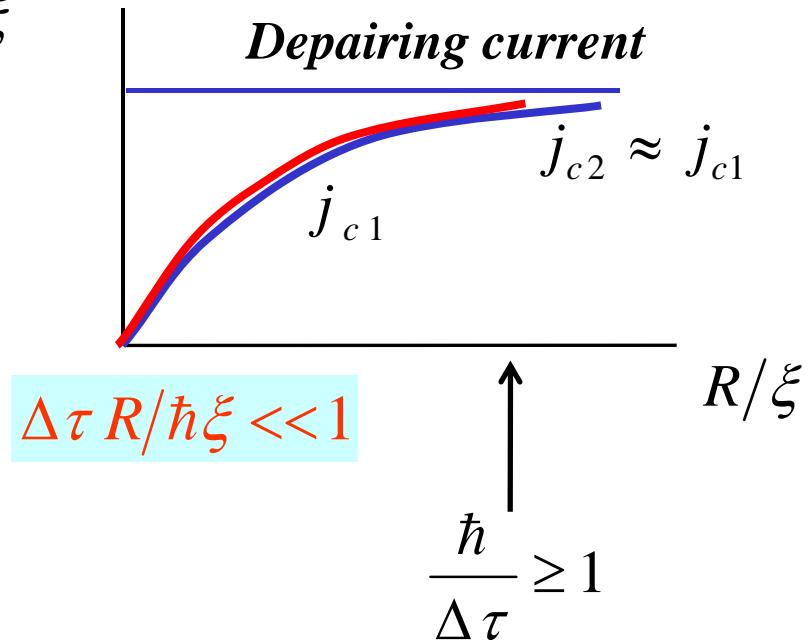
Clean limit



Dirty limit

(in agreement with Thuneberg et al)

$$j_{c1} = j_{depair} \frac{R}{\xi}$$



- Topological electronic transitions in vortex matter
- de Haas – van Alfen oscillations in vortex matter
- Fermi surface transformations and depinning
- bound vortex configurations near the columnar defects