

The new life of the integrability

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Integrability

Classical mechanics, XIX century, Hamilton, Liouville, Lie etc

Phase space Γ_{2n} , ξ^α , $\alpha = 1, \dots, 2n$, $\{\xi^\alpha, \xi^\beta\} = \Omega^{\alpha\beta}(\xi)$

Poisson structure, $\{f(\xi), g(\xi)\} = \Omega^{\alpha\beta} \frac{\partial f}{\partial \xi^\alpha} \frac{\partial g}{\partial \xi^\beta}$

$\Omega^{\alpha\beta} = -\Omega^{\beta\alpha}$, Yacobi identity

Evolution: H energy

$$\frac{d}{dt} \xi^\alpha = \{H, \xi^\alpha\} = -\Omega^{\alpha\beta} \frac{\partial H}{\partial \xi^\beta}$$

Complete integrability: $H + (n-1)$ integrals of motion Q_i in involution

$$\{H, Q_i\} = 0, \quad \{Q_i, Q_j\} = 0$$

Angle-Action variables

$$\xi \rightarrow (I^k, \theta^k), \quad k = 1, \dots, n$$

$$\{I^i, I^k\} = 0, \quad \{\theta^i, \theta^k\} = 0, \quad \{I^i, \theta^k\} = \delta^{ik}$$

$$\begin{cases} \frac{d}{dt} I^k = 0 \\ \frac{d\theta^k}{dt} = \phi^k(I) \end{cases}$$

Korteweg de Vries equation

$$v_t = 6vv_x + v_{xxx}, \quad -\infty < x < \infty, \quad v(x) \rightarrow 0, |x| \rightarrow \infty$$

$$\{v(x), v(y)\} = \delta'(x - y), \quad H = \int (v^3 + \frac{1}{2}v_x^2) dx$$

$$\text{soliton} \quad v(x, t) = \frac{A}{\cosh^2 a(x - vt)}$$

Method of inverse scattering

Gardner, Green, Kruskal, Miura 1967

$$\psi'' + k^2\psi = v(x)\psi, \quad \psi \rightarrow \begin{cases} e^{ikx}, & x \rightarrow -\infty \\ t(k)e^{ikx} + r(k)e^{-ikx}, & x \rightarrow \infty \end{cases}$$

$$|t(k)|^2 + |r(k)|^2 = 1$$

$$v(x) \leftrightarrow \{r(k), \lambda_k, c_k\}$$

$$v(x) \longrightarrow v(x, t)$$



$$r(k) \longrightarrow r(k, t)$$



$$r(k, t) = e^{ik^3t}r(k, 0)$$

Complete integrability

Zakharov, Faddeev 1971

Action variables

$$\rho(k) = \frac{\sqrt{\lambda}}{2} \ln(1 - |r(\sqrt{\lambda})|^2), \lambda_k$$

Angle variables

$$\theta(k) = \arg r(k), \theta_k$$

$$P = \sum \lambda_i + \int_0^\infty \lambda \rho(\lambda) d\lambda$$

$$H = \sum -(-\lambda_i)^{3/2} + \int_0^\infty \lambda^{3/2} \rho(\lambda) d\lambda$$

solitons \Leftrightarrow particles

Further development

70-ties, numerous examples, continuous or discrete space variables NS, Toda, SG etc

Finite interval, periodic conditions \Rightarrow elliptic functions

Sine-Gordon

$$\varphi_{tt} - \varphi_{xx} + \frac{m^2}{\gamma} \sin \gamma\varphi = 0$$

Relativistic model

Solitons — new particles

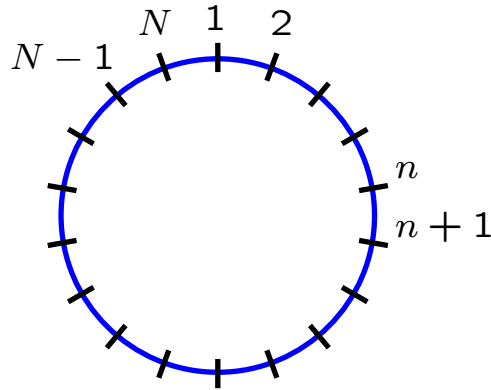
Breathers — their bound states

Duality weak–strong $\gamma \leftrightarrow \frac{1}{\gamma}$

Quantization, first quasiclassical, then from 1979 quantum

Algebraic Bethe Ansatz: Sklyanin, Takhtajan, Korepin, Izergin, Kulish, Semenov-Tian-Shansky, Reshetikhin

Heisenberg spin 1/2 XXX chain



$$\vec{s}_n = (s_n^1, s_n^2, s_n^3)$$

$$[s_n^a, s_m^b] = i\delta_{nm}\epsilon^{abc}s_n^c$$

$$[s_n^a, s_m^a] = 0 \quad n \neq m$$

$$\text{spin } 1/2, \quad \mathbb{C}^2, \quad s^a = \frac{1}{2}\sigma^a$$

$$\mathfrak{H} = \prod \otimes^N \mathbb{C}^2, \quad s_n^a = I \otimes I \dots s^a \dots I$$

$$H = J \sum_{n=1}^N \vec{s}_n \vec{s}_{n+1}, \quad \vec{s}_{N+1} = \vec{s}_1$$

Ferromagnet $J < 0$

Degenerate ground state,

$$E_0 = \frac{1}{4} J N$$

$\Omega =$ spins up

$$S^3 = \sum_n s_n^3 \quad S^3 \Omega = \frac{1}{2} N J \Omega$$

$N \rightarrow \infty$ symmetry breaking

Excitations: magnons and bound states

Antiferromagnet $J > 0$

one vacuum; “Dirac sea”

Excitations: spin 1/2 spinons

No symmetry breaking

Instructive examples in QFT

Quantum realization of the Inverse Scattering Method

$$L_n(\lambda) = \begin{pmatrix} \lambda + is_n^3 & is_n^- \\ is_n^+ & \lambda - is_n^3 \end{pmatrix} \quad s_n^\pm = s_n^1 \pm s_n^2$$

Block matrix or 4×4 matrix in $\mathbb{C}^2 \otimes \mathbb{C}^2$

$$L_n(\lambda) \sim \begin{array}{|c|} \hline \color{blue}{+} \\ \hline \end{array}$$

$$\text{FCR: } R(\lambda - \mu)L_n^1(\lambda)L_n^2(\mu) = L_n^2(\lambda)L_n^1(\mu)R(\lambda - \mu)$$

$$R(\lambda) \sim \color{blue}{\times}$$

$$\begin{array}{|c|} \hline \color{blue}{\times} \\ \hline \end{array} = \begin{array}{|c|} \hline \color{red}{\times} \\ \hline \end{array}$$

$$\mathbb{C}^2 \otimes \mathbb{C}^2 \otimes \mathbb{C}^2$$

$$\psi_{n+1} = L_n(\lambda)\psi_n \quad M_N(\lambda) = \overleftarrow{\prod} L_n(\lambda) = \begin{pmatrix} A(\lambda) & B(\lambda) \\ C(\lambda) & D(\lambda) \end{pmatrix}$$

$$R(\lambda - \mu)M^1(\lambda)M^2(\mu) = M^2(\lambda)M^1(\mu)R(\lambda - \mu)$$

$$T_N(\lambda) = \text{tr } M(\lambda) = A(\lambda) + D(\lambda)$$

$$[T_N(\lambda), T_N(\mu)] = 0$$

Generating function for the set of $N - 1$ independent operators

$$H = \frac{dT(\lambda)}{d\lambda} T^{-1}(\lambda)|_{\lambda=i/2}$$

Together with S^3 complete set of commuting conserved quantities

Generalizations

1. Higher spin
2. Anisotropy XYZ

$$H = \sum (J^1 s_n^1 s_{n+1}^1 + J^2 s_n^2 s_{n+1}^2 + J^3 s_n^3 s_{n+1}^3)$$
$$J_1 = J_2, \quad \text{XXZ}$$

3. Other groups and/or representations
 4. Lattice spacing Δ
 5. Inhomogenous chain $L_n(\lambda) \rightarrow L_n(\lambda - \lambda_n)$
- In particular alternating $\lambda_n = -\lambda_{n+1}$

All known integrable models are different limits – universality of spin chains

Higher spin XXZ

Kulish-Reshetikhin

$$[s_n^+, s_n^-] = \frac{\sin \gamma s_3}{\sin \gamma}$$

Quantum groups

FCR history:

1. factorized scattering

Berezin, Yang, Brezin–Zinn-Justin

2. Boltzman weight in 2-dim. model of classical statistical physics
Onzager, Lieb, Baxter

Yang-Baxter relation

Beautiful universal picture

Unification of several apparently different subjects in MPh

2-dim.: spinons, quantum computers, quantum optics

However recently new signatures of integrability in 4-dim QFT

1. **Lipatov** Reggeization (**Feynman**), $SL(2)$ spin - 1

2. **Minahan, Zarembo**, anomalous dimensions in $N = 4$ SYM, connection with Maldacena duality

Lot of people, **Beisert, Staundacher, Frolov, Arutyunyan, Kazakov**, ... see J. Phys A special volume

3. Vacua in supersymmetrical topological FT
Nekrasov–Shatashvily

New life of integrability